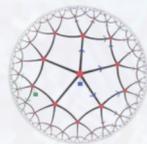


The python's lunch and complexity of decoding Hawking radiation



Caltech



It from Qubit

Simons Collaboration on
Quantum Fields, Gravity and Information

Hrant Gharibyan

IQIM Caltech

Feb 17, 2022

Overview

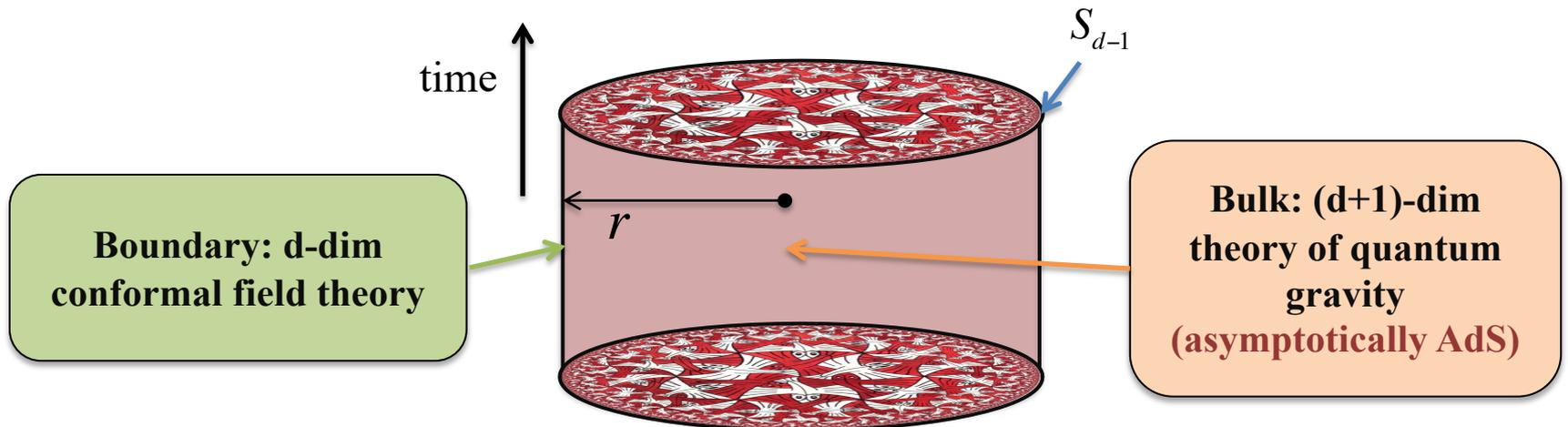
- **Part One:** Introduction to holography and quantum information
- **Part Two:** The python's lunch: geometric obstructions to decoding Hawking radiation

Adam Brown, HG, Geoff Penington, Leonard Susskind; arXiv: 1912.00228

Part One

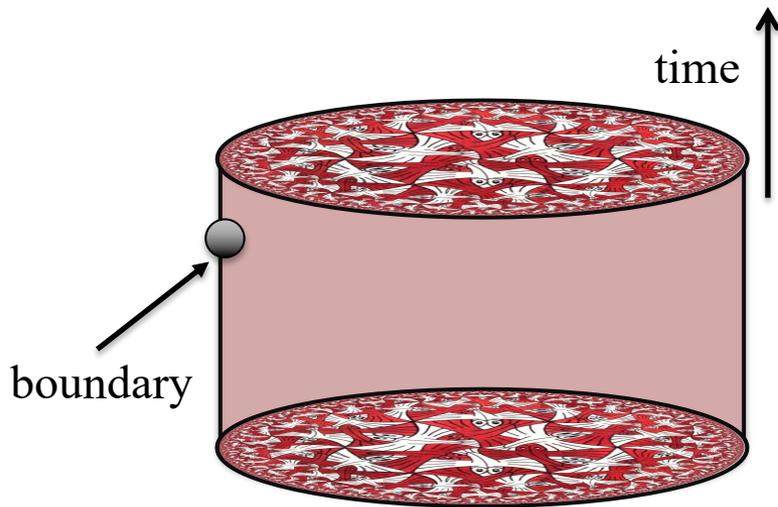
Introduction to holography and quantum information

Holography: AdS/CFT Correspondence



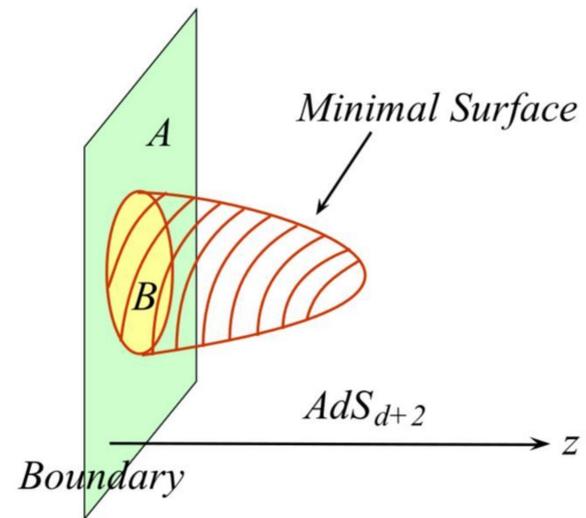
Example: $N=4$ Super-Yang-Mills = Type IIB string theory in $AdS_5 \times S^5$
[Maldecena'97]

Holography and quantum information



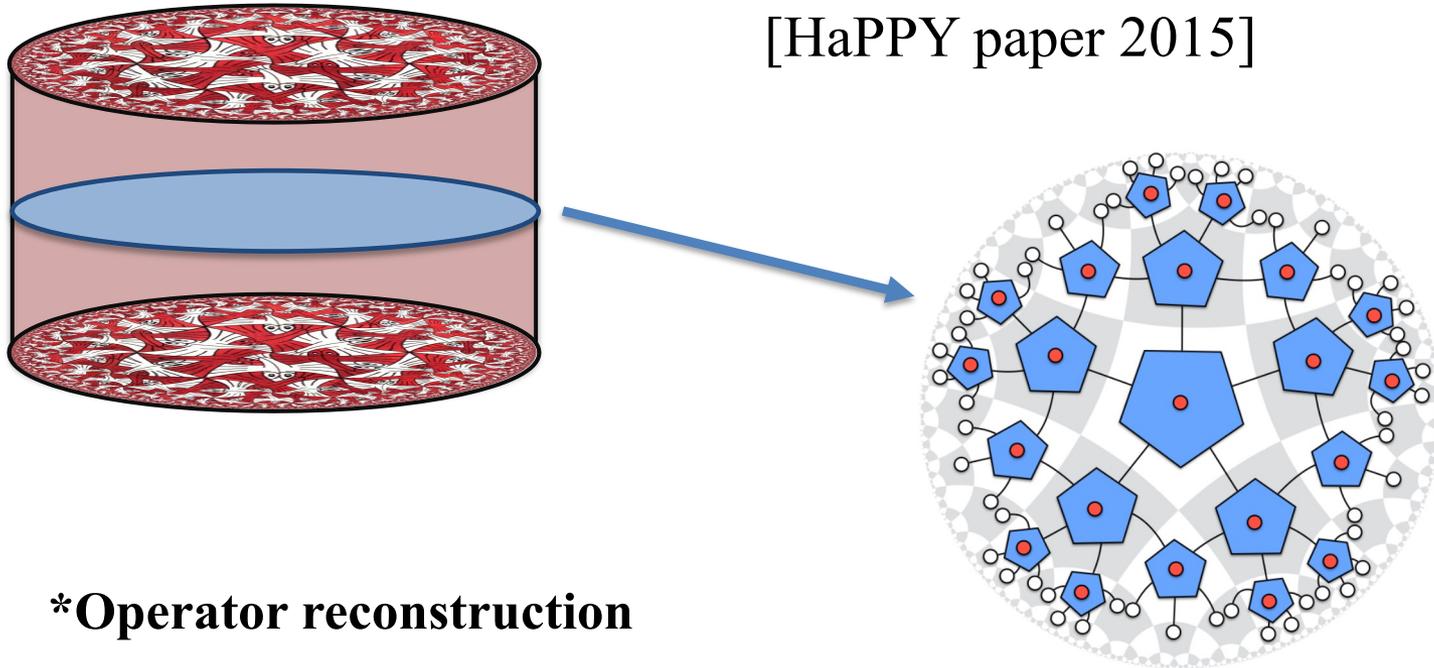
- *there is covariant version
- *as well as generalized entanglement, including bulk entanglement

1. Entanglement entropy = Area [Ryu-Takayanagi 2006]



Holography and quantum information

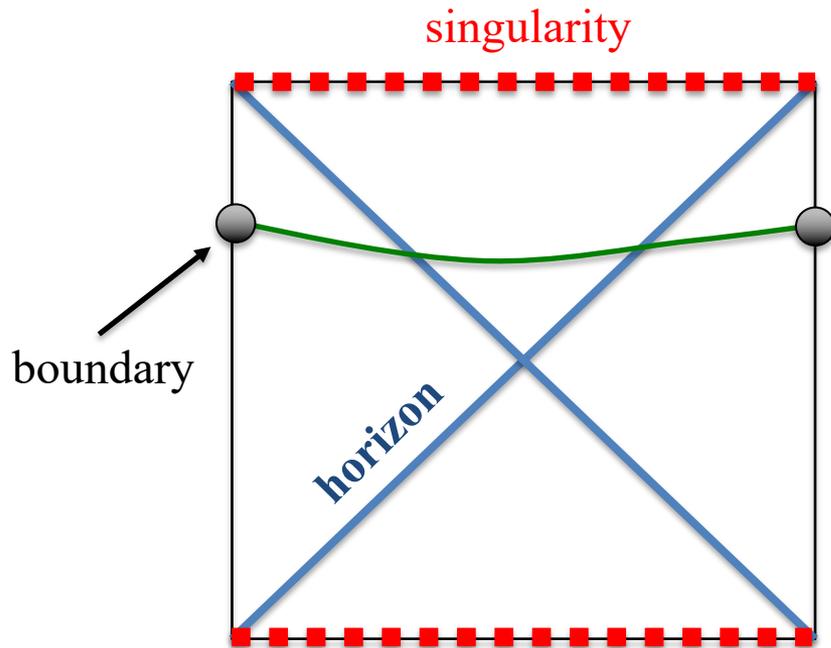
2. Error correction= Space-like slice [HaPPY paper 2015]



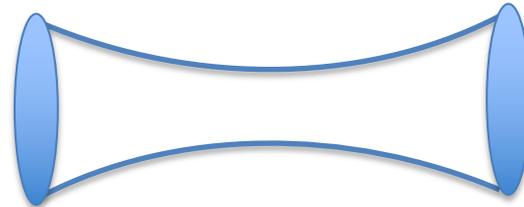
***Operator reconstruction**

***Derivation of the Page curve**

Holography and quantum information

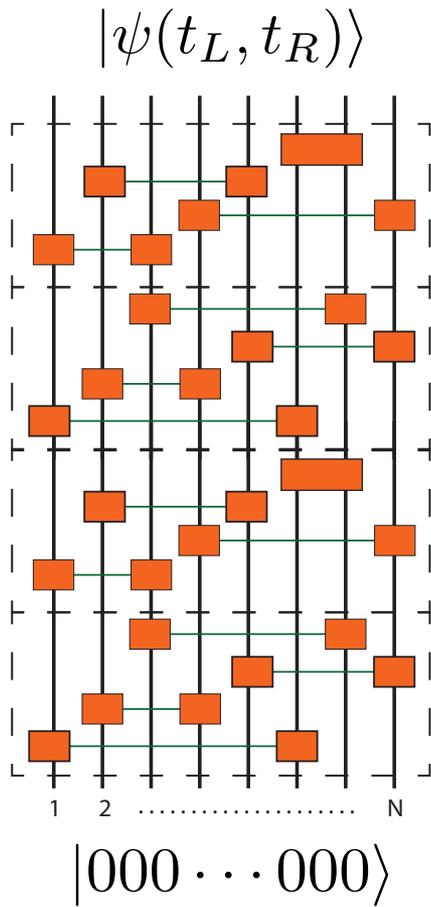


3. Complexity= Volume/Action [Susskind and collaborators 2014]



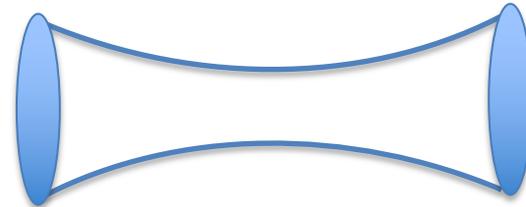
$$\mathcal{C}(|\psi(t_L, t_R)\rangle) \propto \frac{V}{\hbar G l_{AdS}}$$

Holography and quantum information



3. Complexity= Volume/Action

[Susskind and collaborators 2014]



$$\mathcal{C}(|\psi(t_L, t_R)\rangle) \propto \frac{V}{\hbar G l_{AdS}}$$

Part Two

Complexity of evaporating black hole and the python's lunch geometry

[Based on [arXiv: 1912.00228](#)]

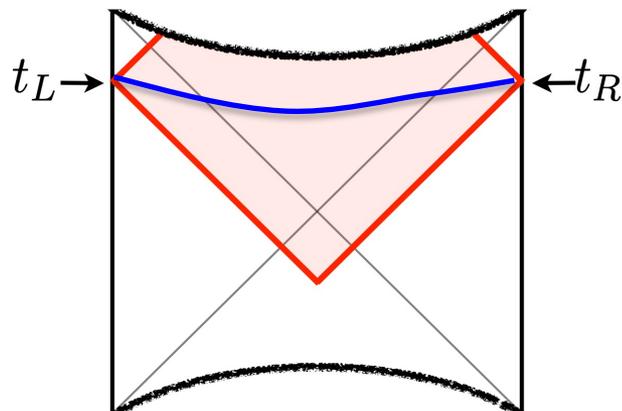
Relationship between complexity and geometry

A. Brown, D. Roberts, L. Susskind, B. Swingle, Y. Zhao

“Complexity Equals Action”; *arXiv: 1509.07876*

D. Stanford, L. Susskind,

“Complexity and Shock Wave Geometries”; *arXiv: 1406.2678*



$$|\psi(t_L, t_R)\rangle$$

State Complexity = Volume of ER Bridge

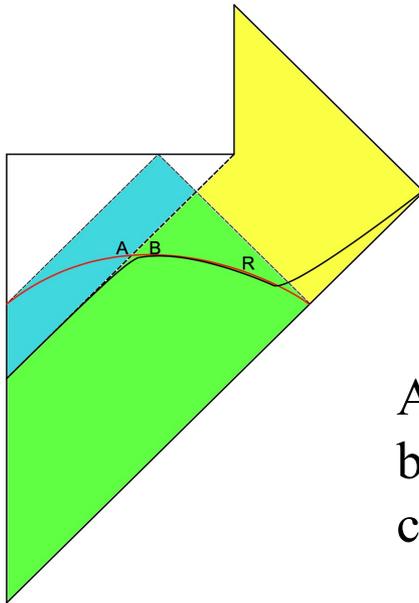
State Complexity = Action of WDW Patch

[Myers and collaborators]

Decoding Hawking radiation is exponentially hard

Daniel Harlow, Patrick Hayden,

“*Quantum Computation vs. Firewalls*”; *arXiv:1301.4504*



Decoding the information that fell into black hole by only action on radiation is hard.

$$\text{Complexity} \sim e^S$$

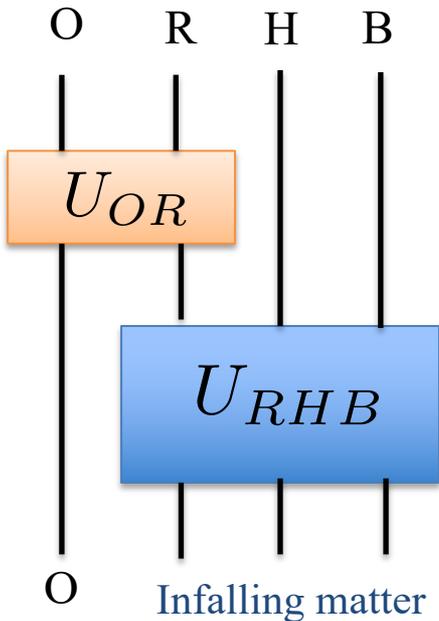
Argument relies on the mixing assumption of black hole and radiation as well as complexity class reduction.

[Harlow-Hayden 2013; Aaronson 2014]

Hawking radiation is pseudorandom (2019)

I. Kim, E. Tang, J. Preskill,

“*The ghost in the radiation: Robust encodings of black hole interior*”; *arXiv:2003.05451*



- RB radiation state is **pseudorandom**
- Computationally bounded observer O cannot destruct encoding of A in RB system.
- Implies that decoding Hawking radiation required exponential complexity

$$\text{complexity} \sim e^S$$

[Kim, Tang, Preskill et. al 2020]

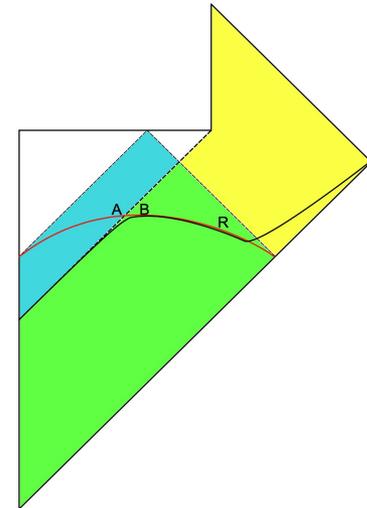
Apparent contradiction

For evaporating black hole volume/action is polynomial and those complexity is polynomial.

poly(S)

Harlow and Hayden suggested it is exponentially hard to decode the radiation.

exp(S)

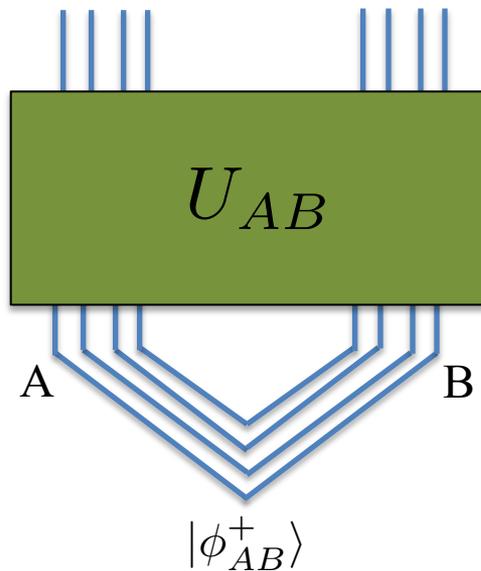


Where is this contradiction coming from?

Restricted vs unrestricted complexity

We are talking about two different notions of complexity

----- unrestricted vs restricted -----



Unrestricted complexity

$$U_{AB} = g_1 g_2 \cdots g_C$$

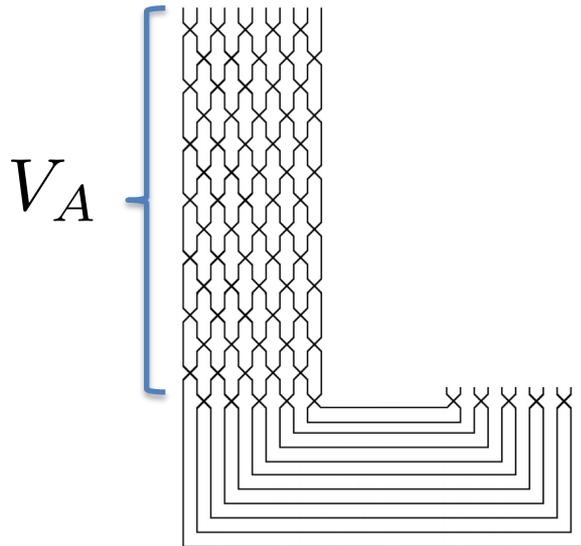
Restricted complexity

$$V_A = g_1 g_2 \cdots g_{C_A}$$

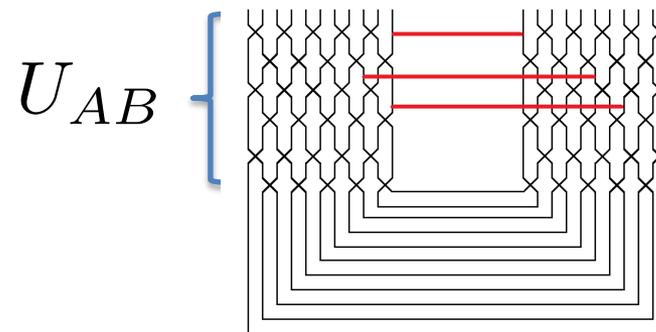
$$V_A |\phi_{AB}^+\rangle \approx U_{AB} |\phi_{AB}^+\rangle$$

Restricted vs unrestricted complexity

Restricted Complexity
(Harlow-Hayden)



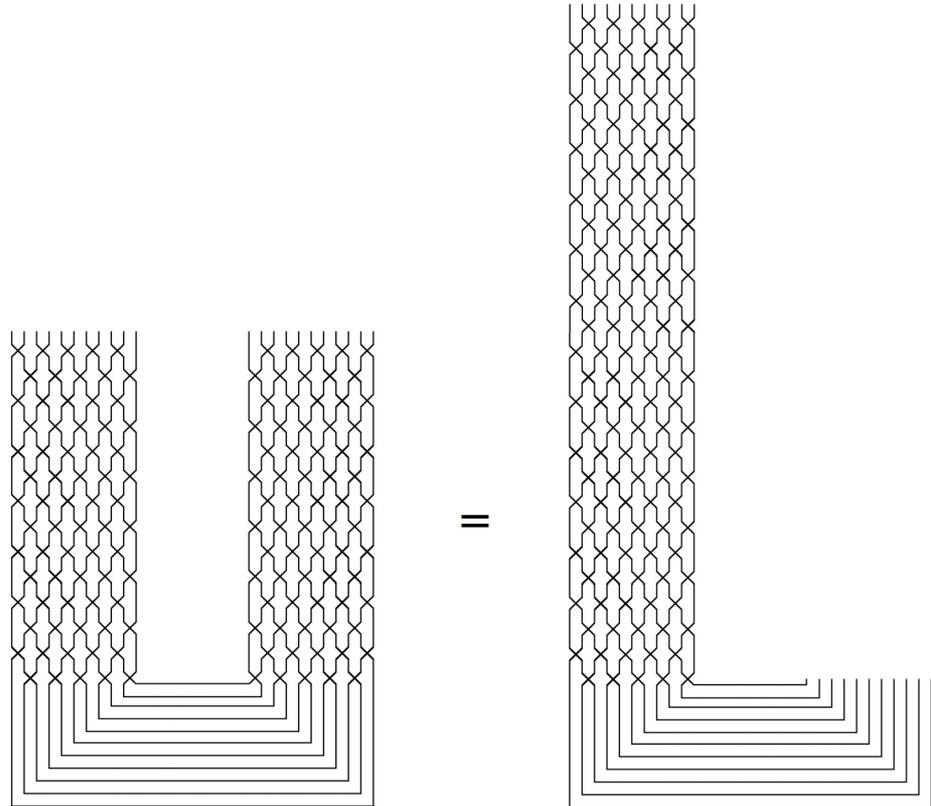
Unrestricted Complexity
(Volume/Action)



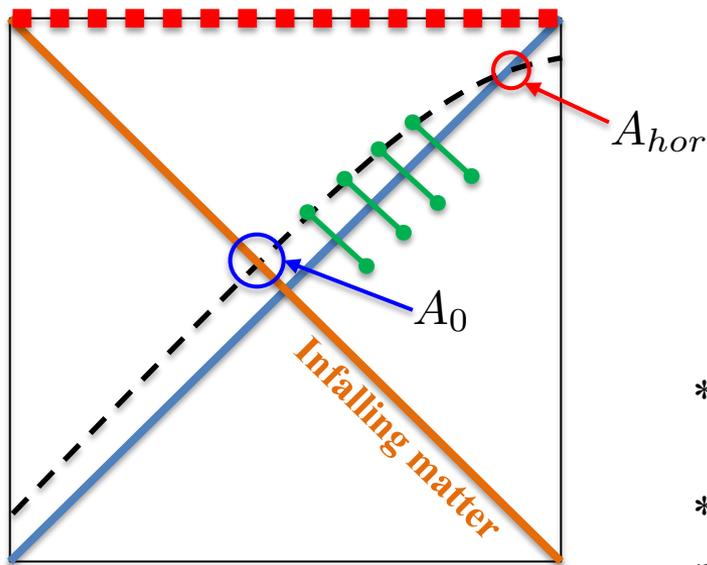
When are restricted and unrestricted complexities very different?

Time evolution of the
thermofield double.

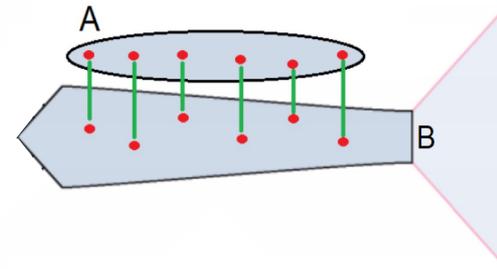
Restricted Complexity
=
Unrestricted Complexity



Cauchy slice of the evaporating black hole



“Nice” Cauchy slice



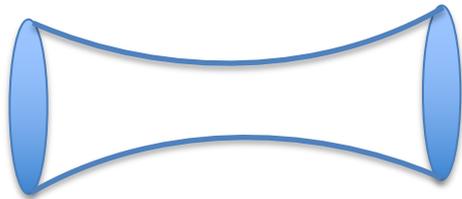
*Radiation and black hole modes are **coupled!**

*A bit after **Page time**, volume/action is polynomial in S .

*Restricted complexity is **exponential**.

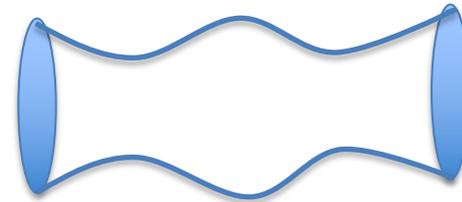
When are restricted and unrestricted complexities very different?

**Small (polynomial)
restricted complexity**



**Time evolved
thermo field double (TFD)**

**Large (exponential)
restricted complexity**



Python's lunch geometry

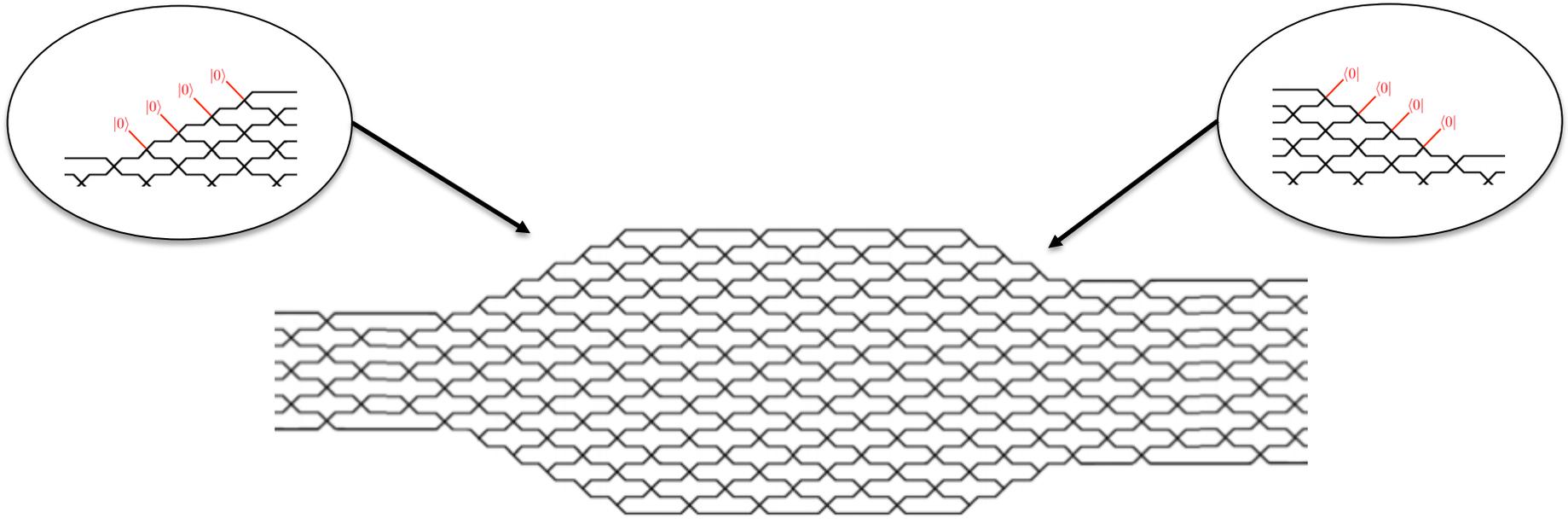
Evaporating Black Hole

When are restricted and unrestricted complexities very different?

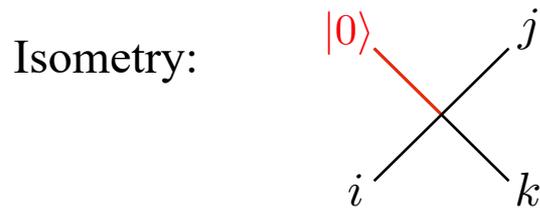
The python's lunch



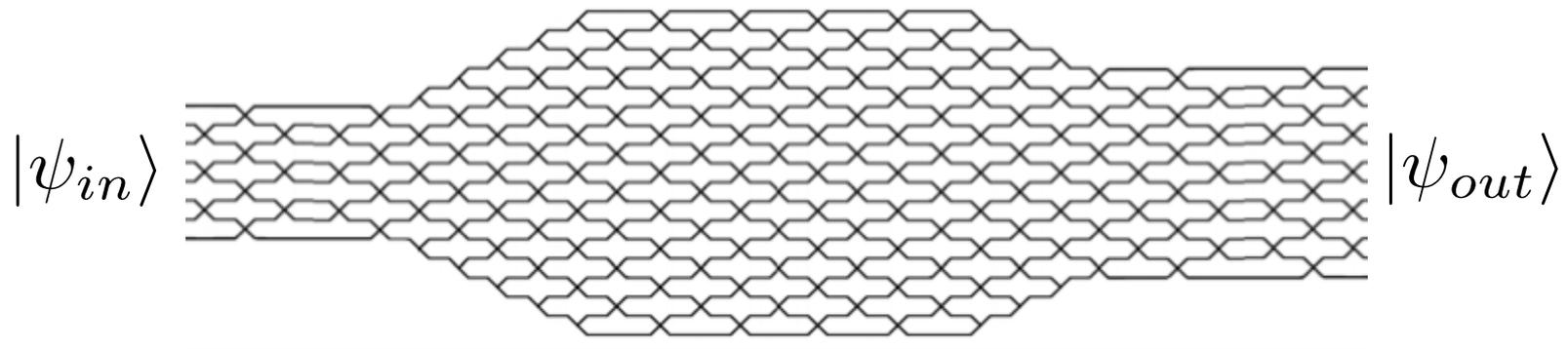
The python's lunch in tensor network (toy model)



Tensor network is built off **unitary gates** and has **isometries** in endpoints.



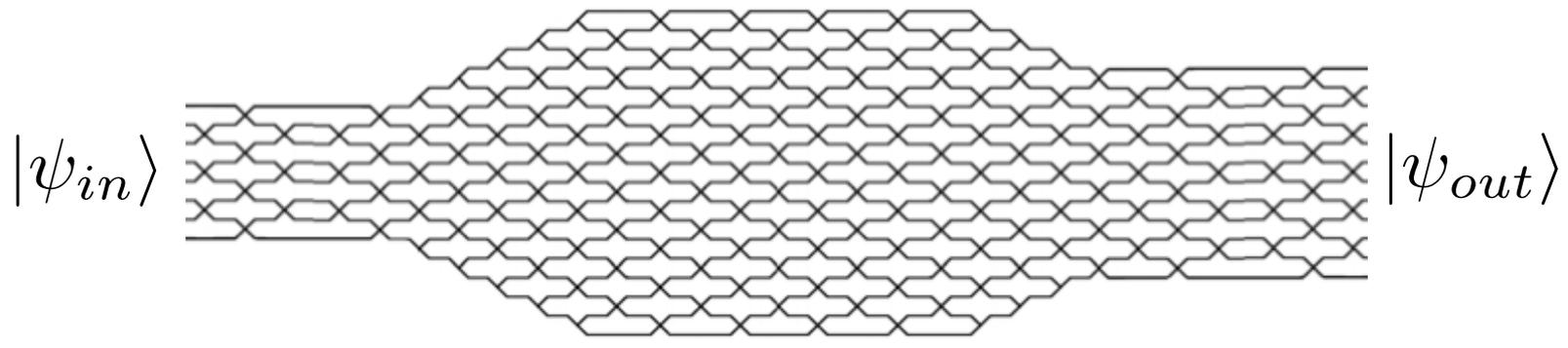
The Python's Lunch in Tensor Network



$$|\psi_{out}\rangle \propto \langle 0|^{m_R} U_{TN} |\psi_{in}\rangle |0\rangle^{m_L}$$

**Ancillary
Input Qubits**

The Python's Lunch in Tensor Network

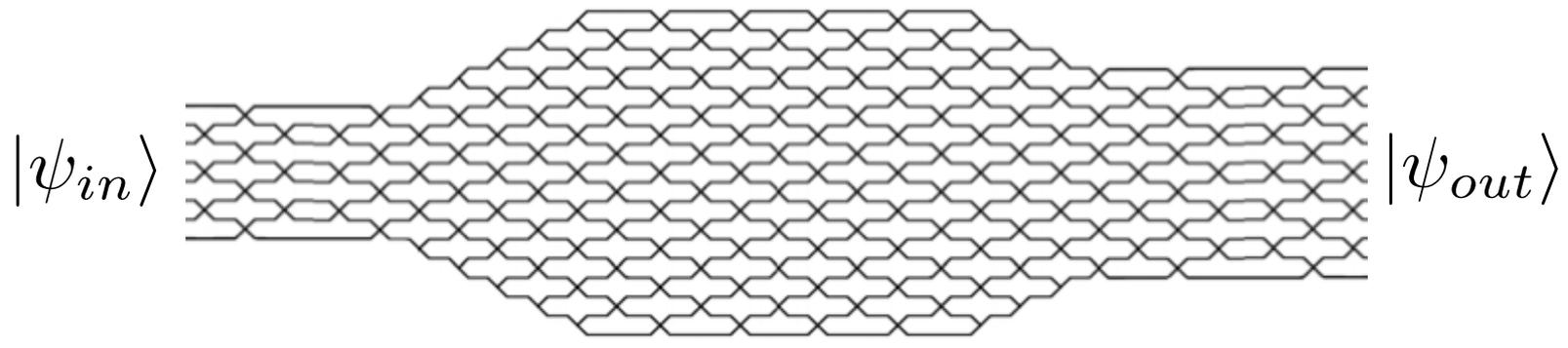


$$|\psi_{out}\rangle \propto \langle 0|^{m_R} U_{TN} |\psi_{in}\rangle |0\rangle^{m_L}$$

**Polynomial
Complexity
Unitary**

**Ancillary
Input Qubits**

The Python's Lunch in Tensor Network



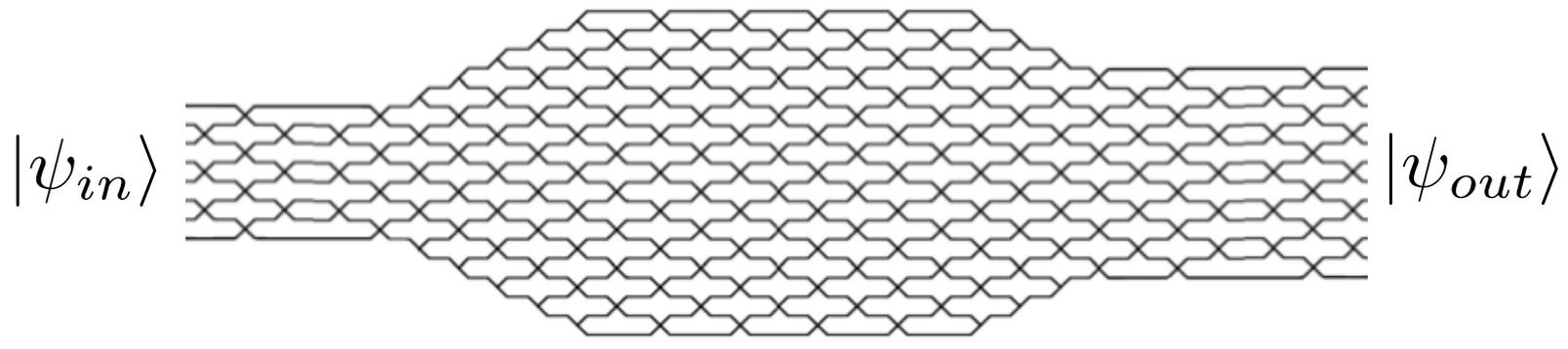
$$|\psi_{out}\rangle \propto \langle 0|^{m_R} U_{TN} |\psi_{in}\rangle |0\rangle^{m_L}$$

Post-selection

**Polynomial
Complexity
Unitary**

**Ancillary
Input Qubits**

The python's lunch: no post-selection

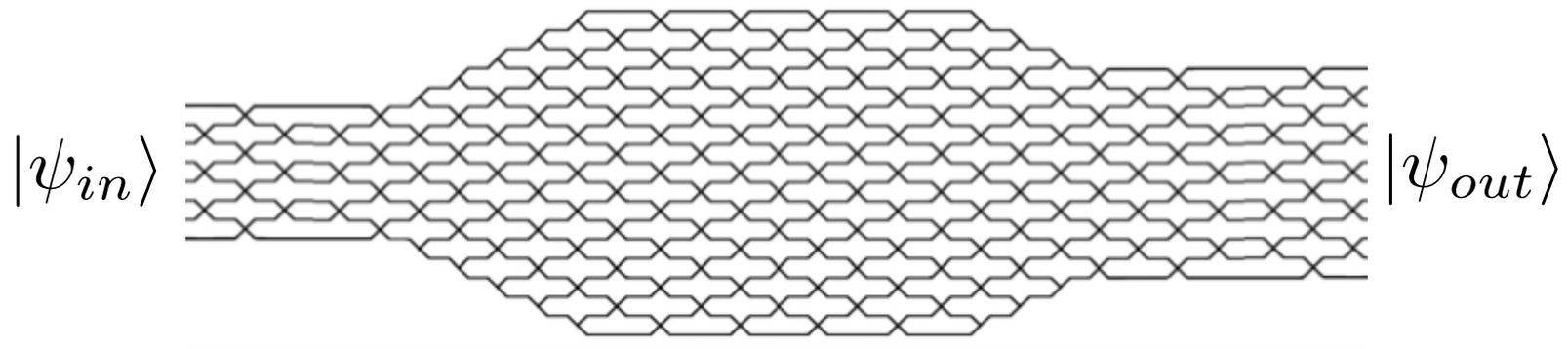


$$|\psi_{out}\rangle|0\rangle^{m_R} = U_{PL}|\psi_{in}\rangle|0\rangle^{m_L}$$

Unitary

**Ancillary
Input Qubits**

The python's lunch: no post-selection



$$|\psi_{out}\rangle|0\rangle^{m_R} = U_{PL}|\psi_{in}\rangle|0\rangle^{m_L}$$

How complex is U_{PL} ?

Unitary

**Ancillary
Input Qubits**

What's the complexity of removing post-selection / final state projection?

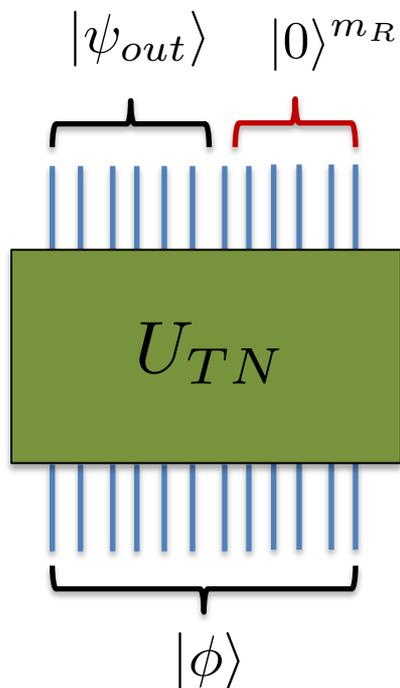
We can measure the qubits and hope to get the right answers. Probability of the right answer is $1/2^m$ and expected complexity

$$C \propto 2^{m_R} \cdot C_{TN}$$

It is not a fixed unitary and complexity can be improved.

We can do a generalized **Grover-like search**.

Grover-like search: state dependent

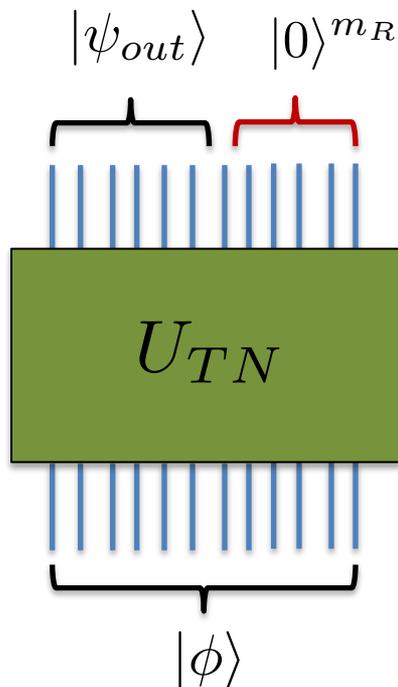


$$U_{\phi} = 2|\phi\rangle\langle\phi| - I$$

$$V = U_{TN}(2|0\rangle\langle 0|^{m_R} - I)U_{TN}^{\dagger}$$

[Kitaev - Yoshida 2017]

Grover-like search: state dependent



$$U_{\phi} = 2|\phi\rangle\langle\phi| - I$$

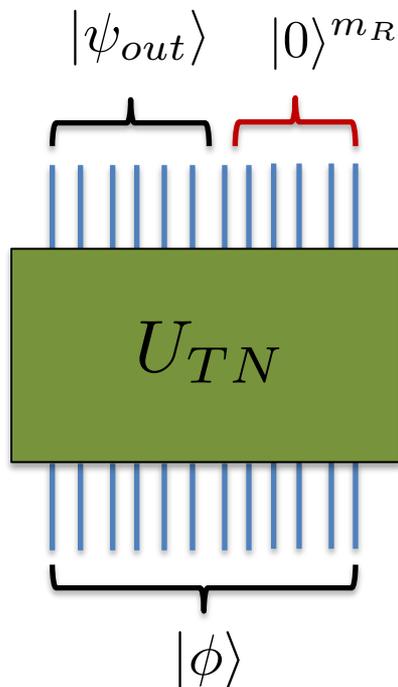
$$V = U_{TN}(2|0\rangle\langle 0|^{m_R} - I)U_{TN}^{\dagger}$$

Repeat sequence: $U_{\phi}V$

$$l = \sqrt{2^{m_R}} \sim \mathbf{times}$$

[Kitaev - Yoshida 2017]

Grover-like search: state dependent



[Kitaev - Yoshida 2017]

$$U_\phi = 2|\phi\rangle\langle\phi| - I$$

$$V = U_{TN}(2|0\rangle\langle 0|^{m_R} - I)U_{TN}^\dagger$$

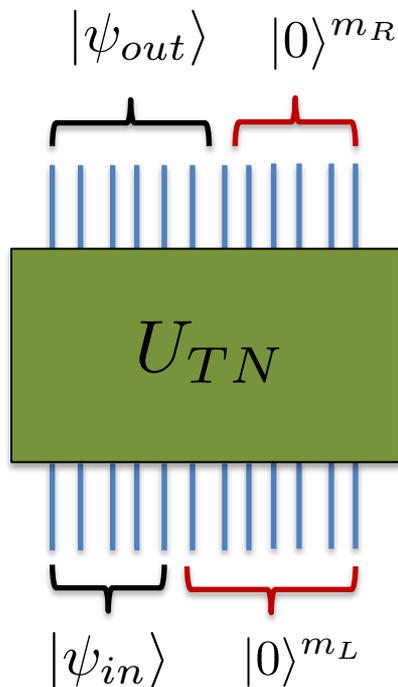
Repeat sequence: $U_\phi V$

$$l = \sqrt{2^{m_R}} \sim \text{times}$$

$$U_{PL} = U_{TN} [U_\phi V]^l$$

This algorithm depends on the state $|\phi\rangle$

Grover-like search: state independent



$$U = 2|0\rangle\langle 0|^{m_L} - I$$

$$V = U_{TN}(2|0\rangle\langle 0|^{m_R} - I)U_{TN}^\dagger$$

$$U_{PL} = U_{TN} [UV]^l$$

$$\mathcal{C}(U_{PL}) \propto \sqrt{2^{m_R}} \cdot \mathcal{C}_{TN}$$

[Berry et. al 2015, Gilyen et. al 2017]

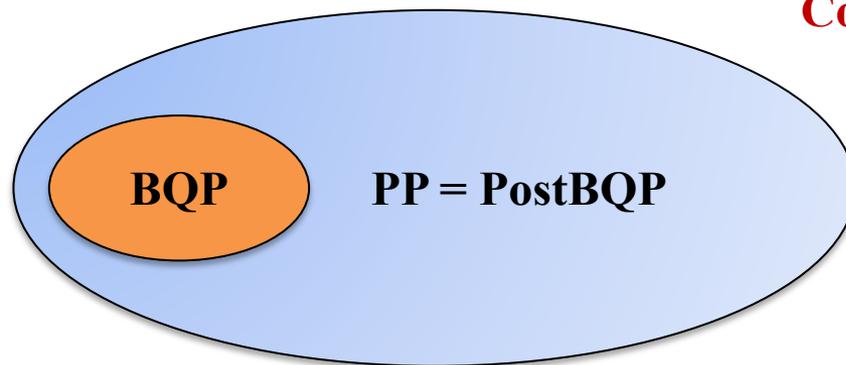
Is Grover search the optimal strategy?

In special cases, for a specific unitary there might be an algorithm that can do better than Grover-like search.

Is Grover search the optimal strategy?

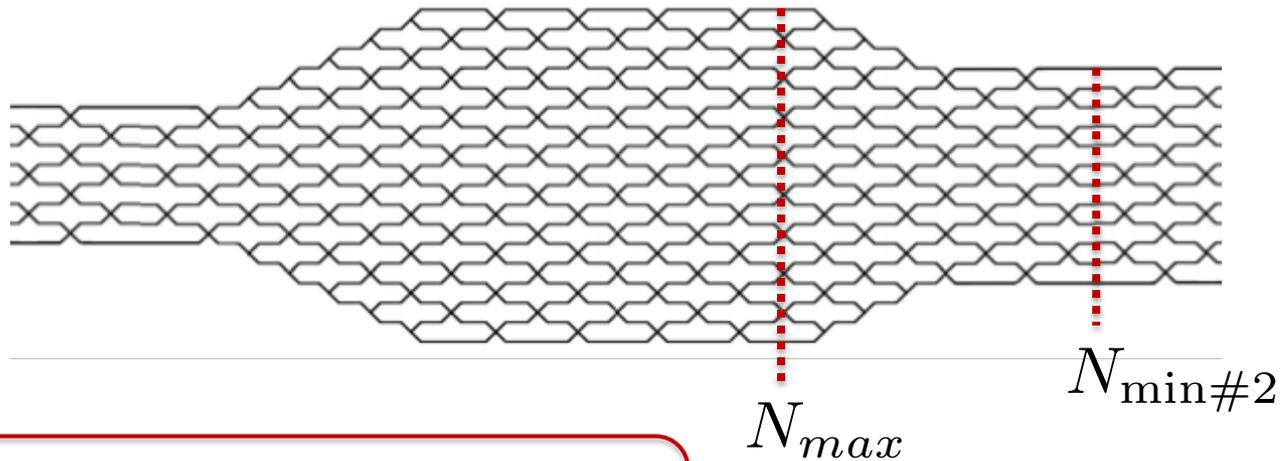
In special cases, for a specific unitary there might be an algorithm that can do better than Grover-like search.

However, for generic, **scrambling tensor network** U_{TN} polynomial algorithm is extremely unlikely to exist.



Complexity class collapse!

Conjecture: restricted complexity in tensor network



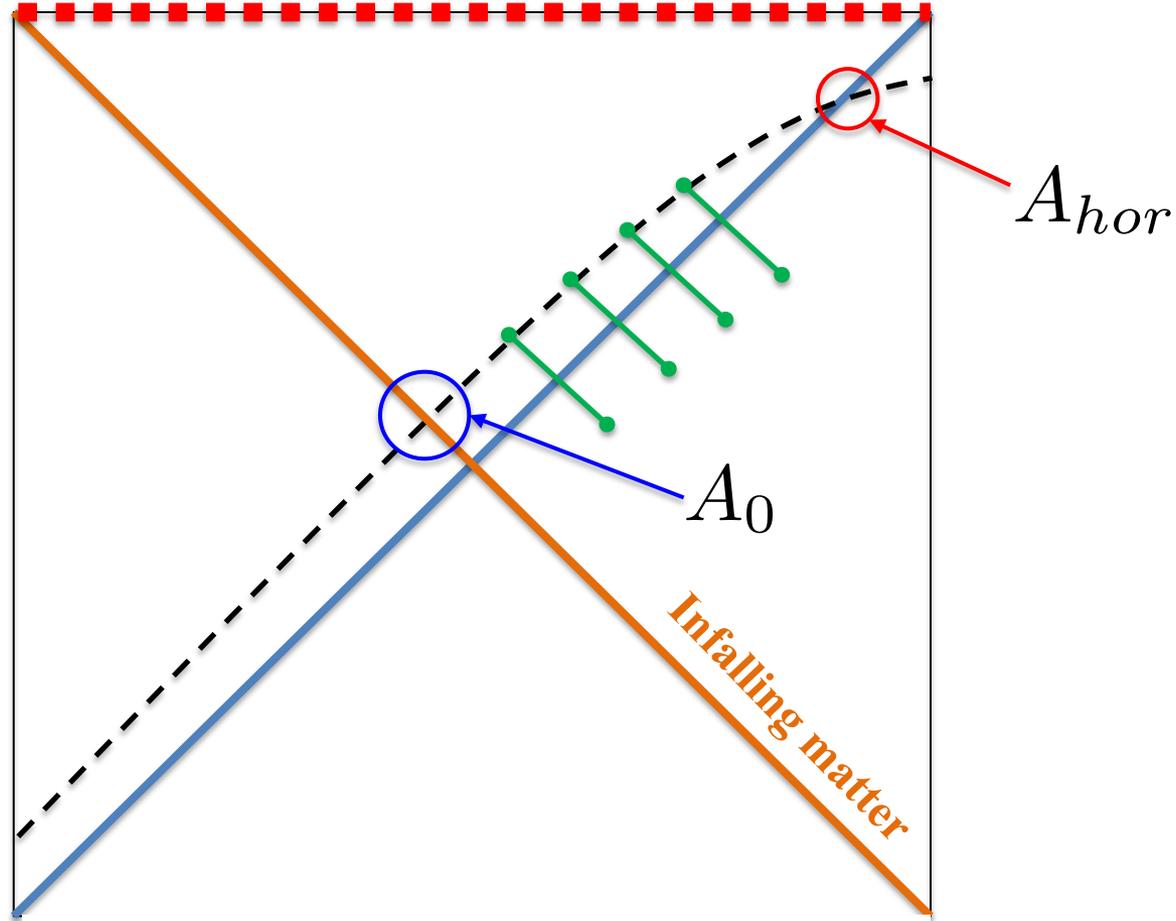
Minimax: minimal of all maximal cuts for all the possible slicing.

Largest minimal cut

The restricted complexity of scrambling tensor network is

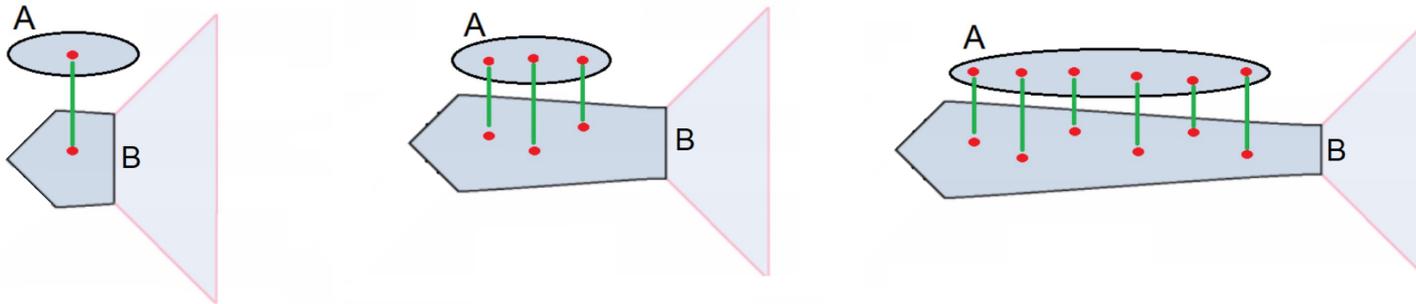
$$\mathcal{C}_R \propto C_{TN} \cdot \exp \left[\frac{1}{2} \left(N_{max} - N_{min\#2} \right) \right]$$

The python's lunch in evaporating black hole



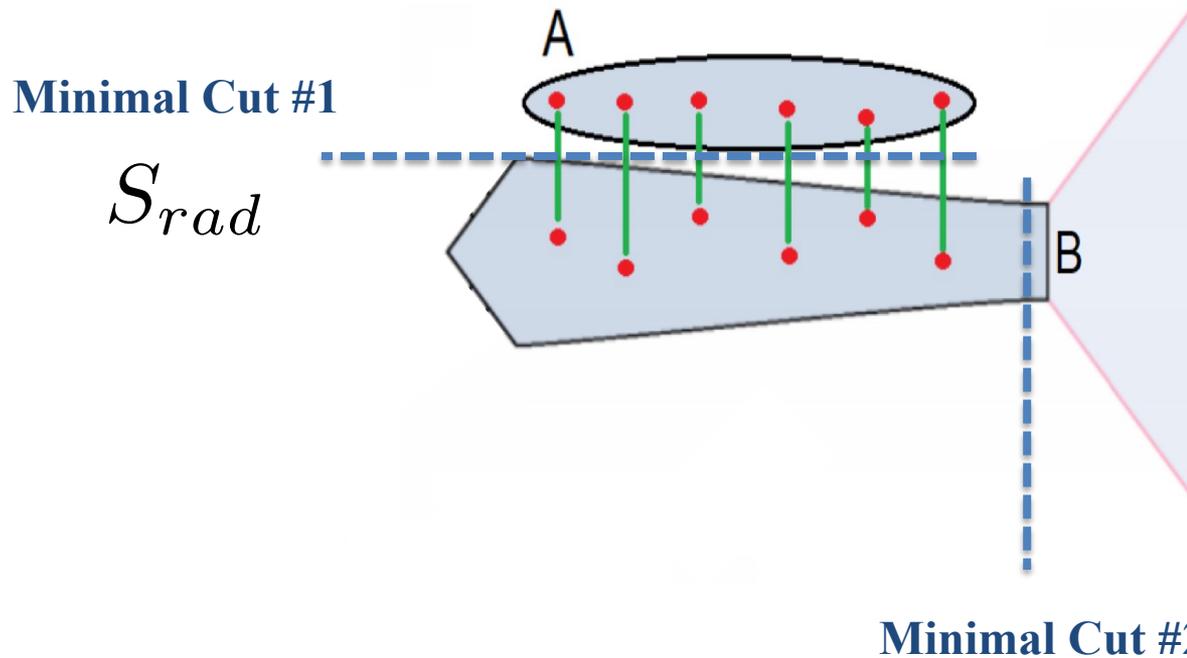
The python's lunch in evaporating black hole

Naïve picture of evaporation process



A is the radiation modes. **B** is the remaining black hole, some nice Cauchy slice going into the horizon.

The python's lunch in evaporating black hole



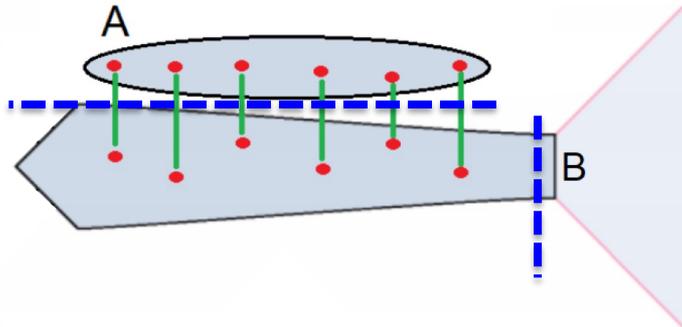
Page time

$$S_{rad} = \frac{A_{hor}}{4\hbar G}$$

What about the maximal cut?

$$\frac{A_{hor}}{4\hbar G}$$

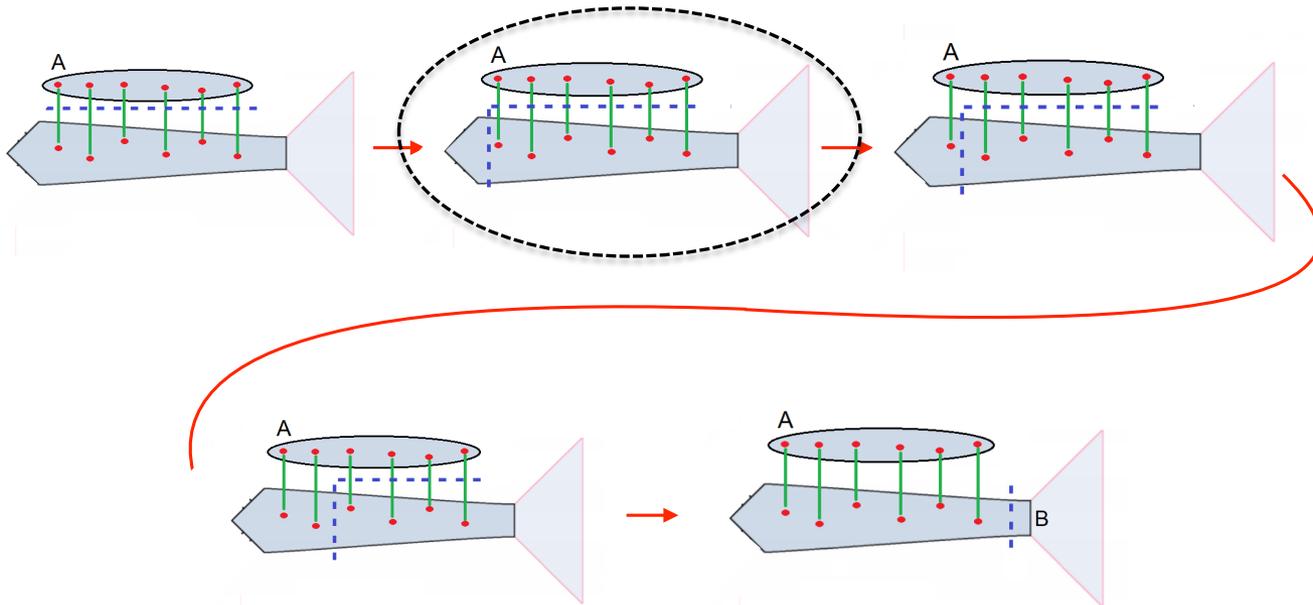
Covariant definition of python's lunch



Both **minimal cuts** are quantum extremal surfaces (extremize the generalized entropy).
[Engelhardt, Wall 2014]

Maximin prescription: find the surfaces with smallest and second smallest generalized entropy then maximize over choice of Cauchy surfaces.
[Wall 2012, AEPU 2019]

Python's lunch search for maximal cut

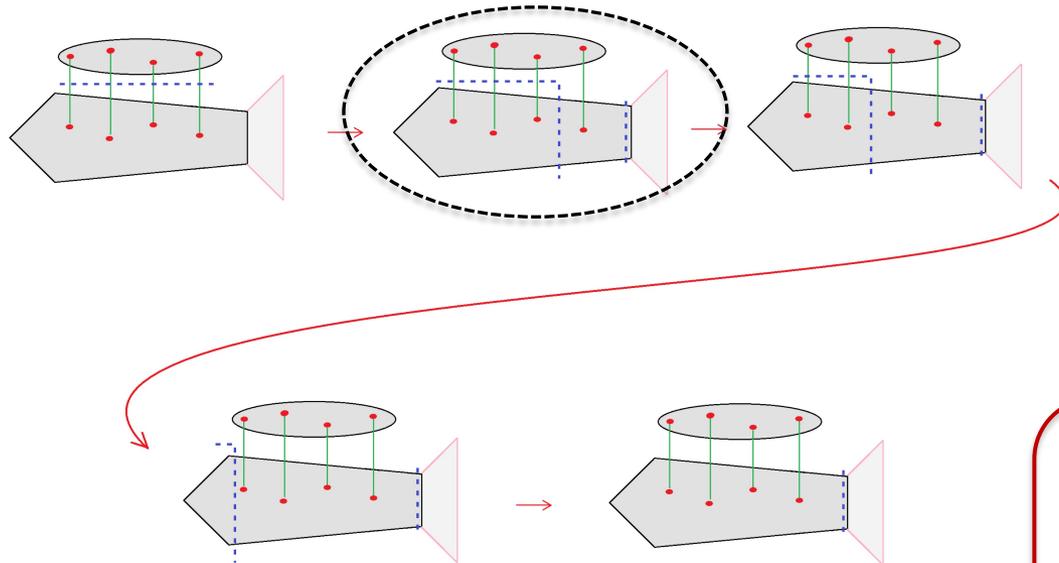


Before half time, we will be to **forward slicing**. The maximal generalized entropy.

$$S_{max}^{(gen)} = S_{rad} + \frac{A_0}{4\hbar G}$$

Minimax: find minimal of maximums of all slices.

Python's lunch search for maximal cut



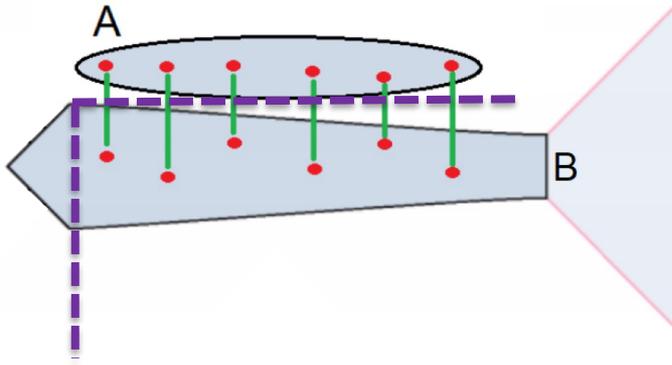
After half time, we will perform **reverse slicing**.

$$S_{max}^{(gen)} = S_{rad} + 2 \frac{A_{hor}}{4\hbar G}$$

Half time

$$A_{hor} = \frac{A_0}{2}$$

Covariant definition of python's lunch



Minimax: maximal cut for each Cauchy slice
find the minimum of the all maximal cuts
for all possible slicing methods. Analogous
to the tensor network story.

Maximinimax: maximize again over all Cauchy surfaces, giving non-minimal
quantum extremal surface.

We find this cuts explicitly for **JT gravity plus non-interacting fermions** and verify the picture from the nice Cauchy slice.

Conjecture: restricted complexity for evaporating black hole

The restricted complexity of evaporating black hole is

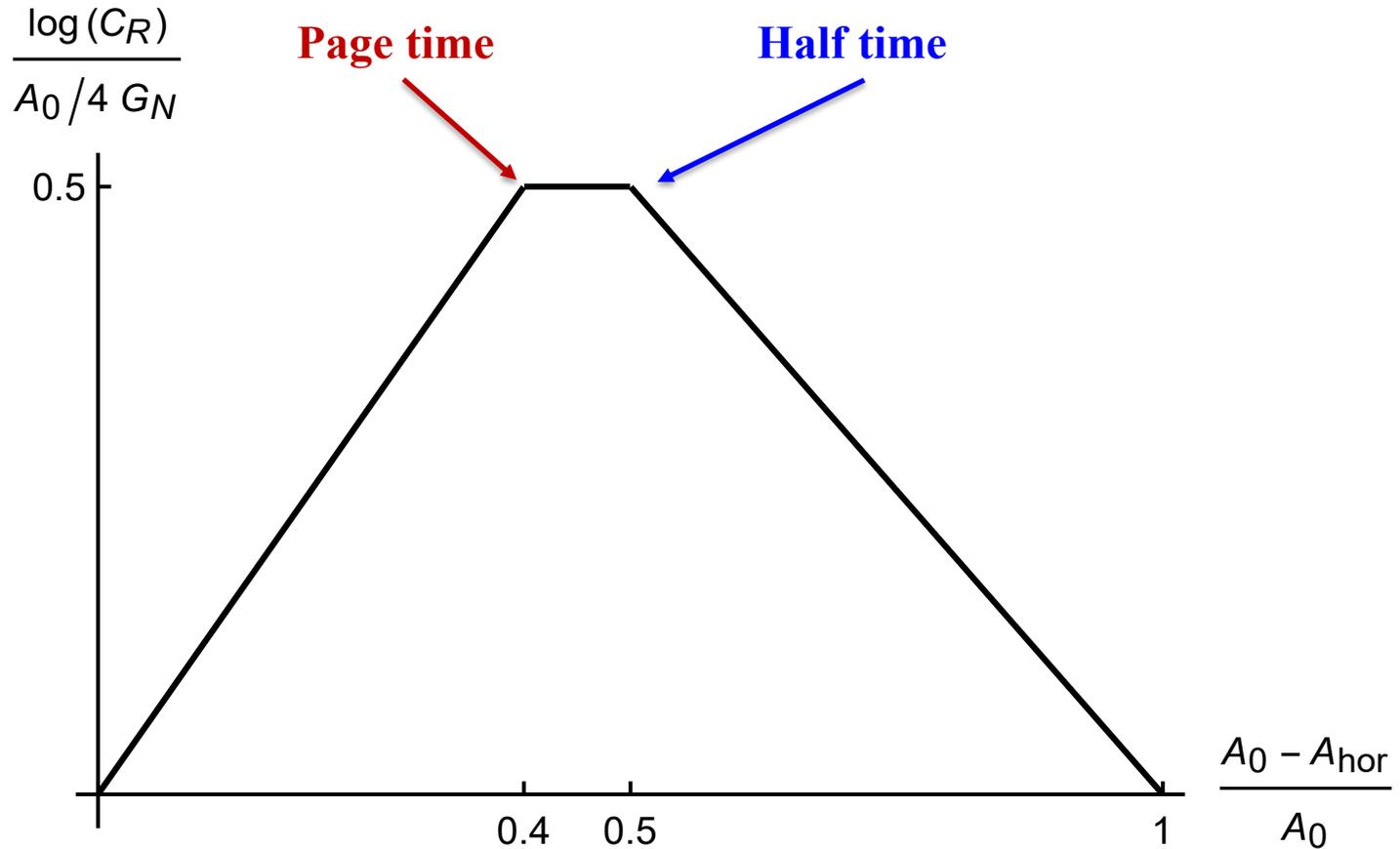
$$C_R \propto C_{TN} \cdot \exp \left[\frac{1}{2} \left(S_{max}^{(gen)} - S_{min}^{(gen)} \right) \right]$$

**Volume/Action
(unrestricted complexity)**

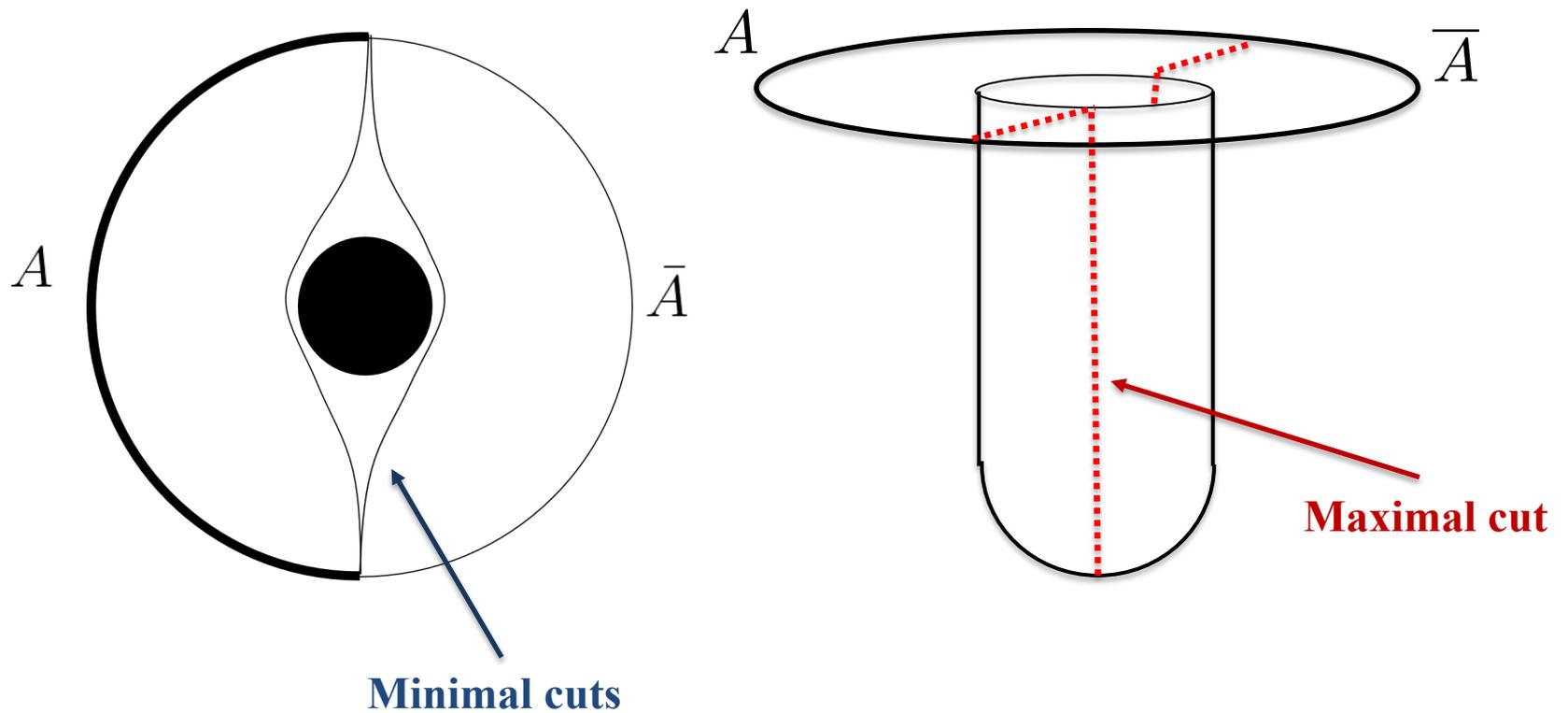
Maximinimax prescription

Maximin prescription

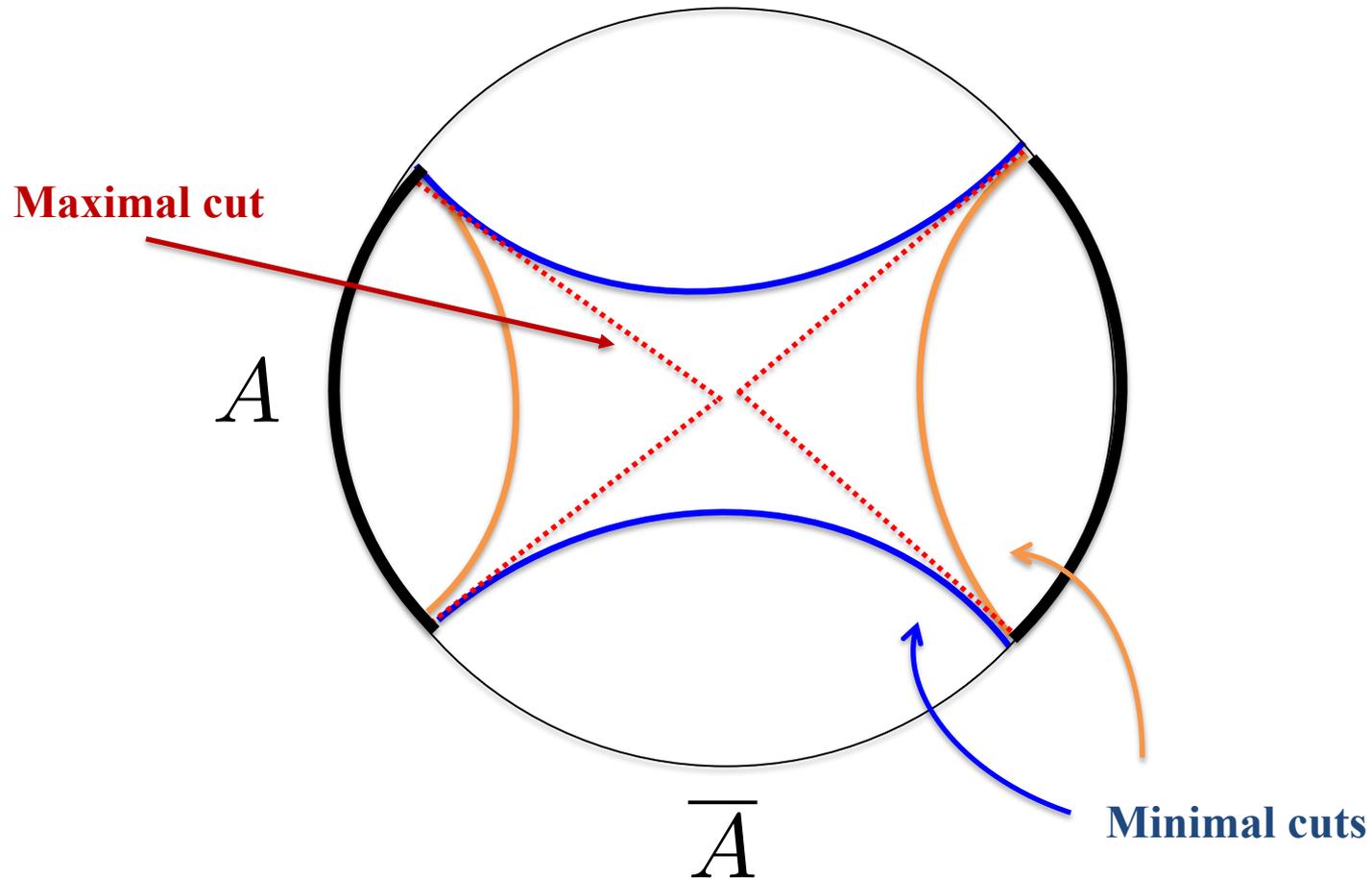
Restricted complexity for evaporating black hole



Python's lunch in pure state black hole

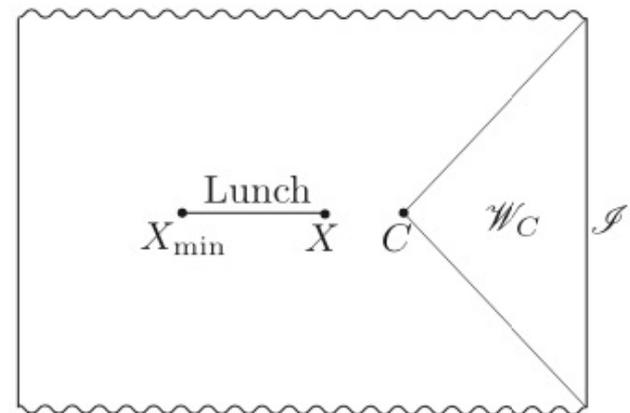
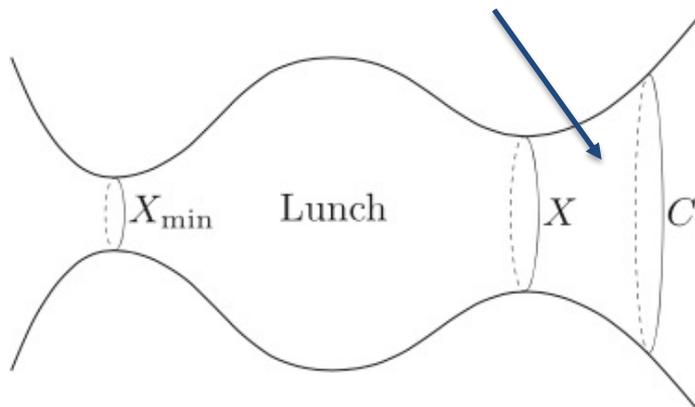


Python's lunch without black holes – empty AdS



Converse of Python's lunch conjecture

Easy to reconstruct



- HKLL construction
- Causal bulk propagation of perturbed boundary condition on Lorenzian timefolds

[Engelhardt, Penington, Shahbazi-Moghaddam 2021]

Final remarks

Summary: The restricted complexity of python's lunch is

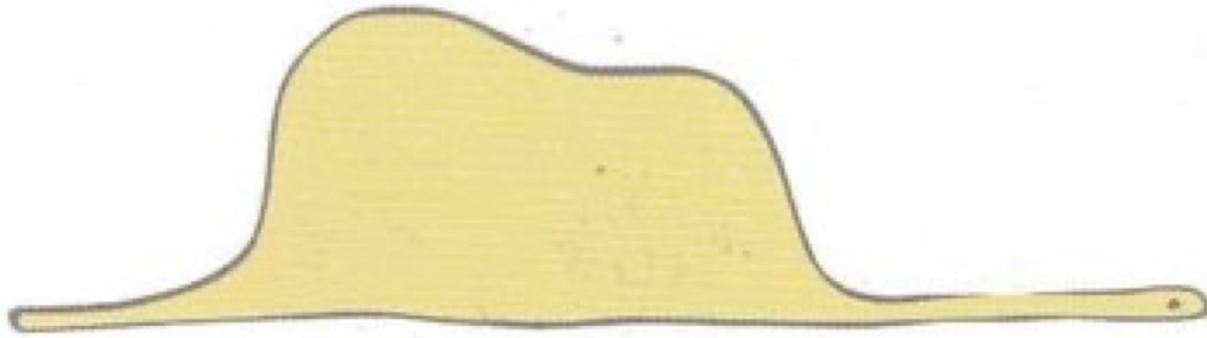
$$C_R \propto C_{TN} \cdot \exp \left[\frac{1}{2} \left(S_{max}^{(gen)} - S_{min}^{(gen)} \right) \right]$$

Future directions

- Test this conjecture in quantum field theories, perhaps with alternative definition of complexity
- Study other examples when python's lunches and non-minimal quantum extremal surfaces appear; pure state black holes and empty AdS.
 - What does this imply about reconstruction of operators in the python's lunch?

The End

The python's lunch



Appendix

Part Three

Quantum gravity in the lab: teleportation by size and traversable wormholes

Based on arXiv: 1911.06314 and 2102.01064

Simulating quantum gravity via holography

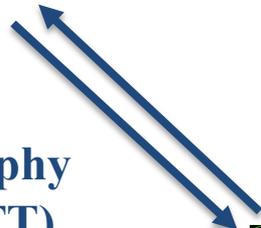
Quantum Devices



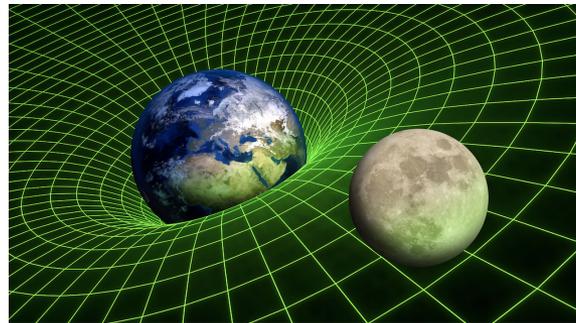
Classical Supercomputers



**Holography
(AdS/CFT)**



Quantum Gravity



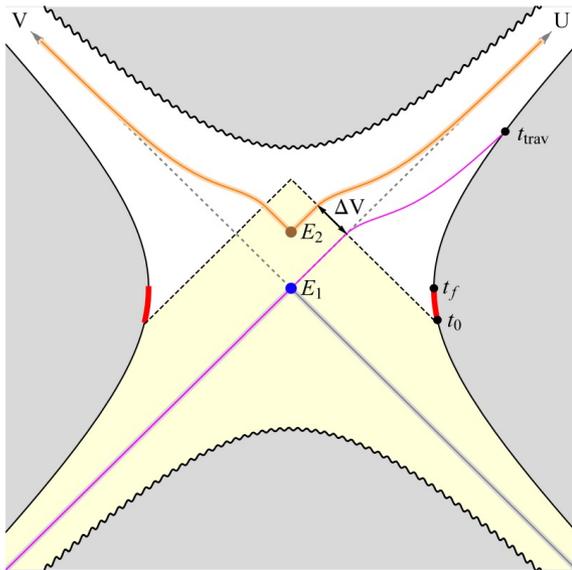
Traversable wormholes in AdS/CFT

P. Gao, D. L. Jafferis, A. Wall

“*Traversable Wormholes via a Double Trace Deformation*”; *arXiv: 1608.05687*

J. Maldacena, D. Stanford,

“*Diving into Traversable Wormholes*”; *arXiv: 1704.05333*

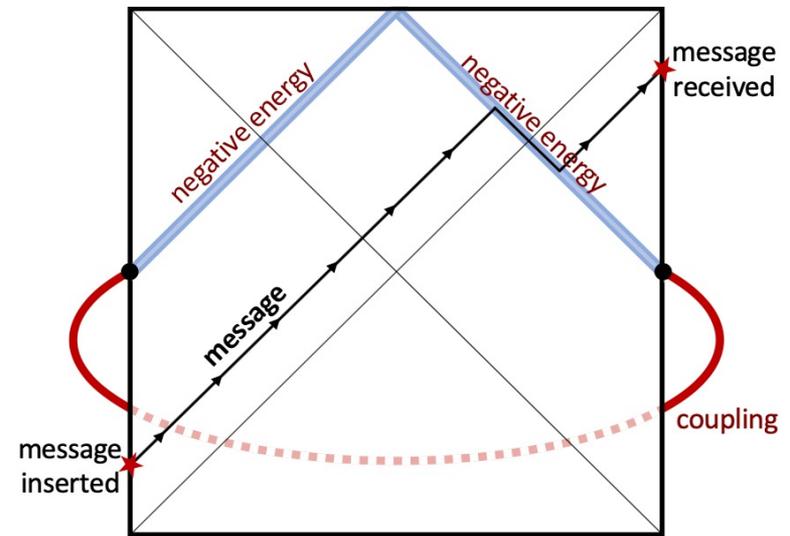
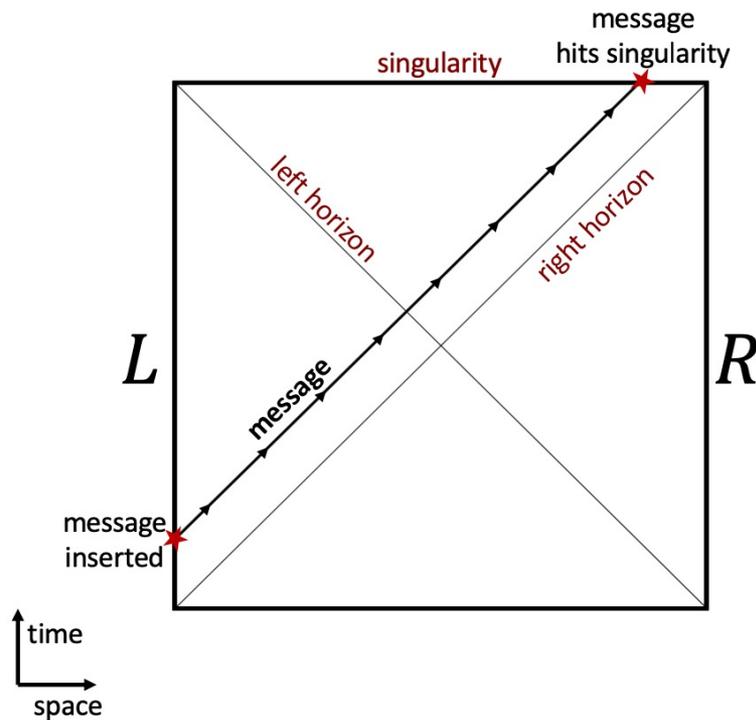


Couple left and right CFT's

$$\exp \left[g \sum_i \mathcal{O}_i^L \mathcal{O}_i^R \right]$$

Sends a negative energy particles in black hole in AdS that makes the wormhole traversable.

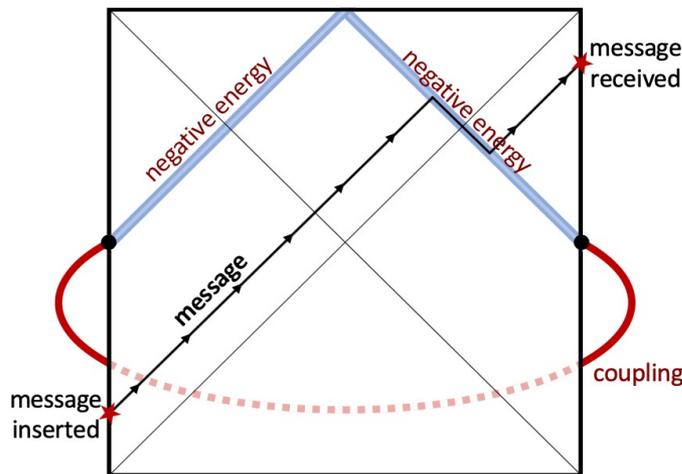
Is traveling through the wormhole a new kind of teleportation protocol?



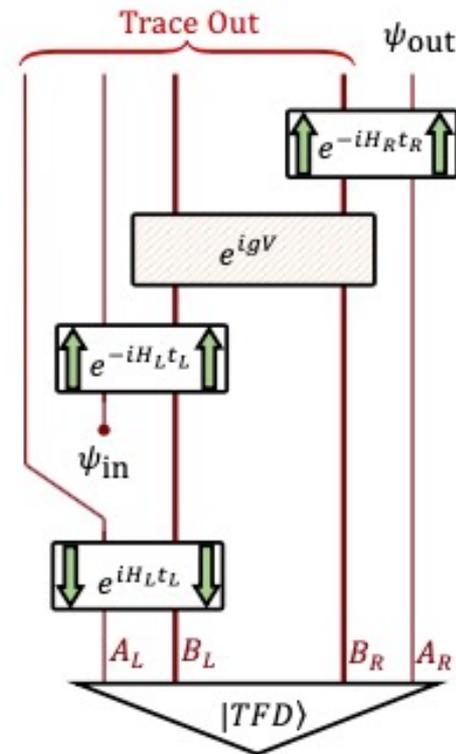
Couple: $\exp \left[g \sum_i \mathcal{O}_i^L \mathcal{O}_i^R \right]$

Understanding the flow of quantum information flow on the boundary

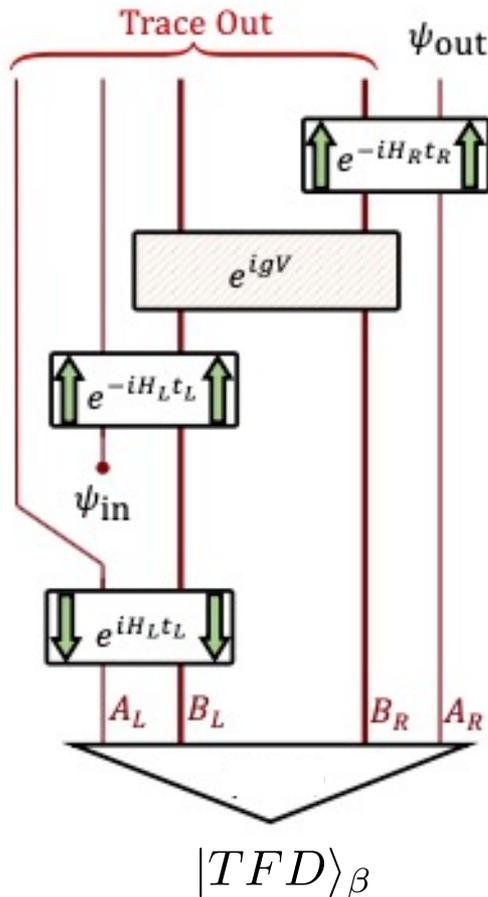
AdS/CFT set up



Boundary interpretation?



Understanding the flow of quantum information flow on the boundary



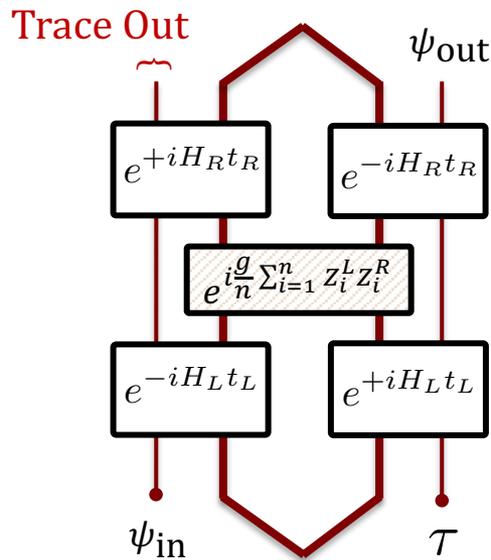
H_L, H_R : left and right Hamiltonians

Couples left-right

$$V = \sum_i O_i^L O_i^R$$

$$|TFD\rangle_\beta = \frac{1}{\sqrt{\text{Tr} e^{-\beta H}}} \sum_a e^{-\beta E_a/2} |E_a\rangle |E_a\rangle$$

Teleportation by size



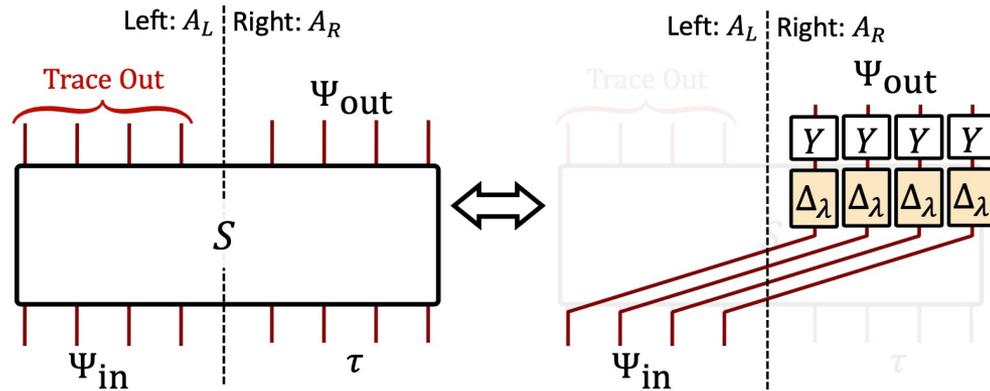
$$\approx \hat{S}$$

**Simple transformation
of the initial circuit**

**Approximately implements
size operator \hat{S} : random matrix,
chaotic spin chains, SYK model**

$$\hat{S}|P\rangle = e^{ig|P|}|P\rangle$$

Teleportation by size



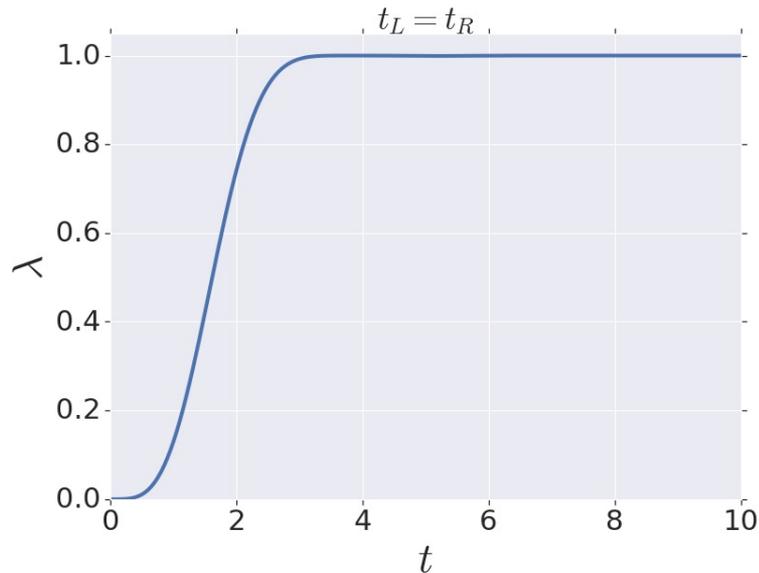
$$\hat{S}|P\rangle = e^{ig|P|}|P\rangle, \text{ where } |P\rangle = P|\phi_{Bell}^+\rangle^{\otimes m}$$

IIXXIIXYIIZIZ ← Pauli strings

|P| is the size of the Pauli operator (number of non-identity elements)

Regime #1: teleportation of a single qubit

1. Teleports a **single qubit** with perfect fidelity.

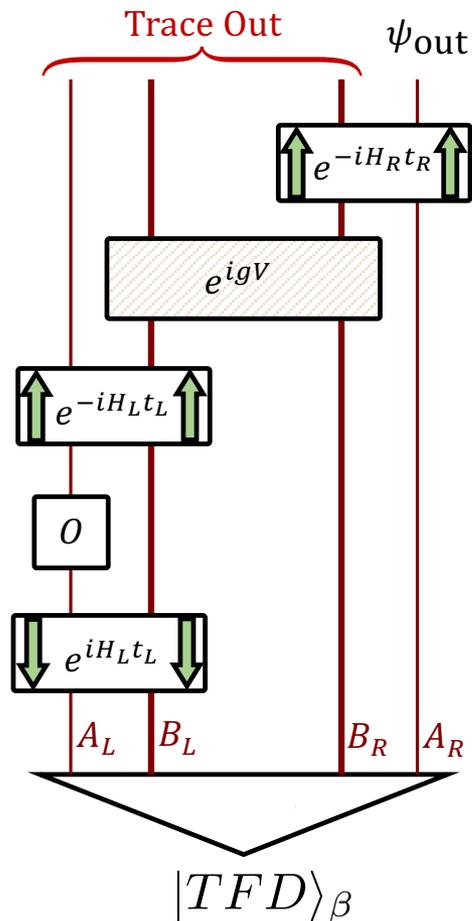


2. Works for **all times** after scrambling time.

3. Works the best at **infinite temperature**.

$$|TFD\rangle_{\beta=0} = |\phi_{Bell}^+\rangle^{\otimes n}$$

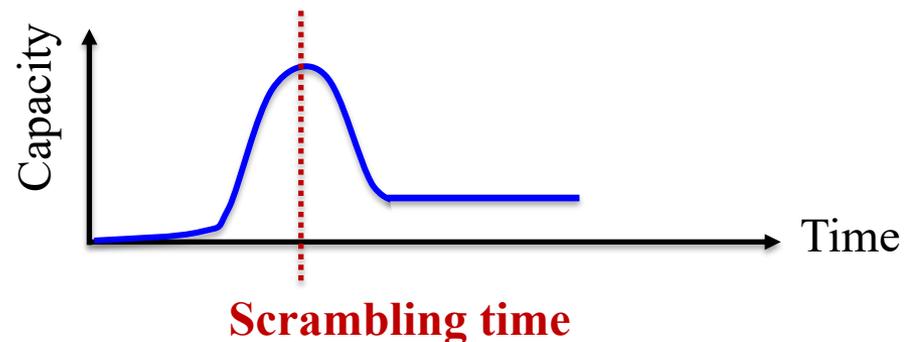
Regime #2: more qubits and low temperatures



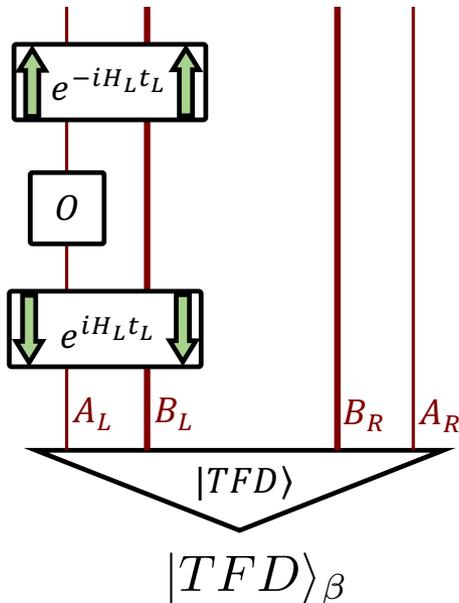
At **low temperature** a new phenomena appears

We can teleport **more than a single qubit**

This works for Hamiltonians with
holographic duals: example is SYK model



Regime #2: size winding explained



Operator growth in thermal ensemble

$$\mathcal{O}(t)\rho_{\beta}^{1/2} = e^{-iHt}\mathcal{O}e^{+iH(t+\frac{\beta}{2})} = \sum_P c_P(t)P$$

complex

Size-winding phenomena

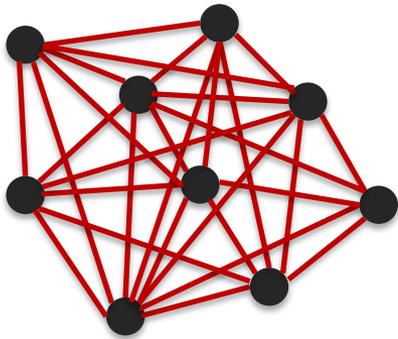
$$q(l) = \sum_{|P|=l} c_P^2 = e^{i\theta} \sum_{|P|=l} |c_P|^2$$

The phase of the **size-distribution winds** near scrambling time:

$$\theta \sim l \times (\text{const})$$

Regime #2: size winding in SYK model

Size-winding is hard to illustrate from the boundary **analytically**.



Sachdev-Ye-Kitaev

$$H_{SYK} = \sum_{i,j,k,l} J_{ijkl} \chi_i \chi_j \chi_k \chi_l$$

We verify the results noticed earlier in JT gravity

Moment = Operator Size

Bulk AdS

Boundary CFT

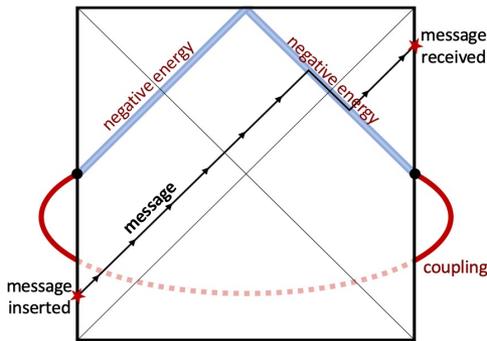
Regime #2: size winding and teleportation

Left operator action

$$\mathcal{O}_L(-t_*)|TFD\rangle_\beta = \sum_P e^{+i\alpha|P|/n} r_P |P\rangle$$

Plus the coupling

$$e^{igV} |P\rangle \approx e^{\frac{i4g|P|}{3n}} |P\rangle$$

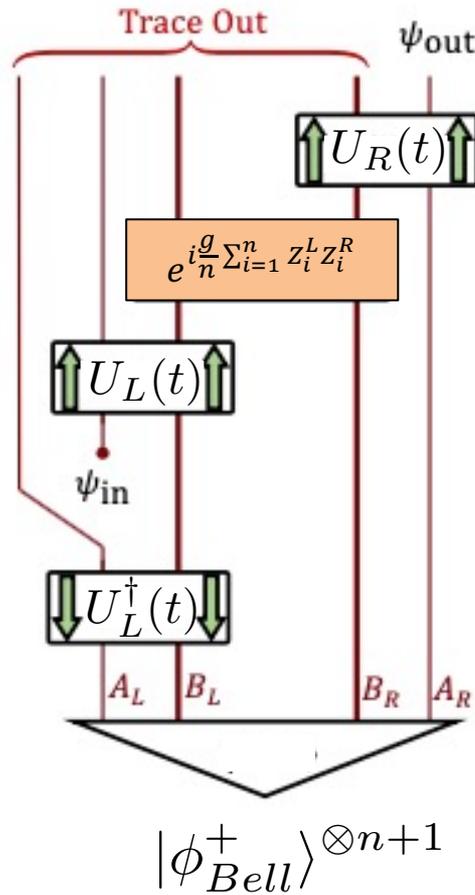


Will equal to an operator acting on the right

$$\mathcal{O}_R(+t_*)|TFD\rangle_\beta = \sum_P e^{-i\alpha|P|/n} r_P |P\rangle$$

[Closely related work by Schuster et.al [arXiv:2102.00010](https://arxiv.org/abs/2102.00010)]

Experimental realizations



We discuss several experimental realizations that are within reach:

Superconducting qubits,
Rydberg atoms, Trapped ions

Floquet version of Rydberg dynamics

$$U_R(t) = U_L(t) = U^t$$

where each time step is $U = U_K U_L$

$$U_K = \exp \left[ib \sum_i X_i \right]$$

$$U_L = \exp \left[i \sum_i h_i Z_i + iJ \sum_i Z_i Z_{i+1} \right]$$

What can we learn about quantum gravity (holography) using the tools of quantum information and computation?

- ❖ Complexity = Volume conjecture for evaporating black hole and **the python's lunch** [[arXiv:1912.00228](#)]

$$C_R \propto C_{TN} \cdot \exp \left[\frac{1}{2} \left(S_{max}^{(gen)} - S_{min}^{(gen)} \right) \right]$$

- ❖ Holographic complexity in **JT gravity** [[arXiv:1810.08741](#)]
- ❖ Onset of chaos in **SFF** from random quantum circuits [[arXiv:1803.08050](#)]

Future research:

- Find other examples of holography where python's lunch appears
- Test this conjecture in quantum field theories, perhaps with alternative definition of complexity

Can we study deep questions about quantum gravity in the near term quantum devices? (quantum gravity in the lab)

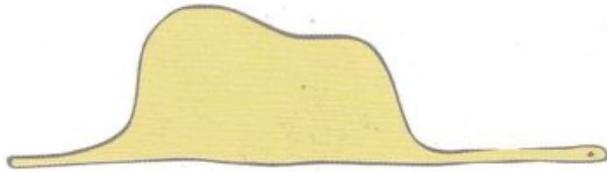
- ❖ **Teleportation by size** protocol inspired by traversable wormholes
- ❖ The **size winding** phenomenon and connection to semi-classical gravity duals [[arXiv:1911.06314](#) and [2102.01064](#)]
- ❖ Simulating **matrix quantum mechanics** on a quantum computer (BFSS/BMN holographic dual to string theory) [[arXiv: 2011.06573](#)]

Future research:

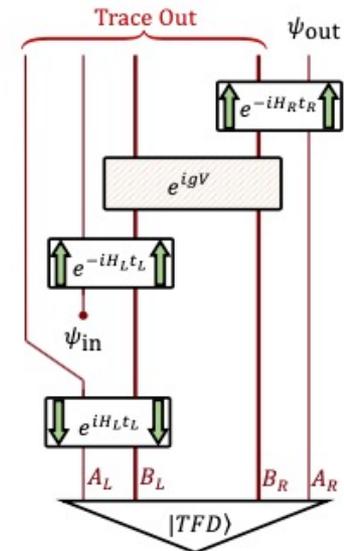
- What is the link between maximal chaos (in OTOC) and size winding?
- Teleportation protocol for GHZ-like entangled state
- Construct new protocols and observables that have quantum gravity interpretation and can be measured in near term experiments

The End

The python's lunch



Quantum gravity in the lab

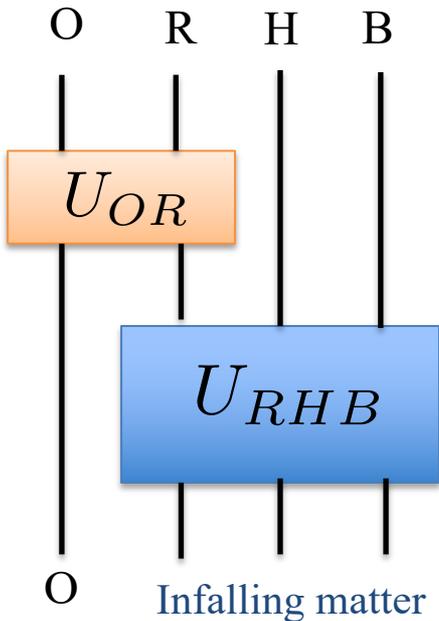


Appendix

Hawking radiation is pseudorandom (2019)

I. Kim, E. Tang, J. Preskill,

“*The ghost in the radiation: Robust encodings of black hole interior*”; *arXiv:2003.05451*

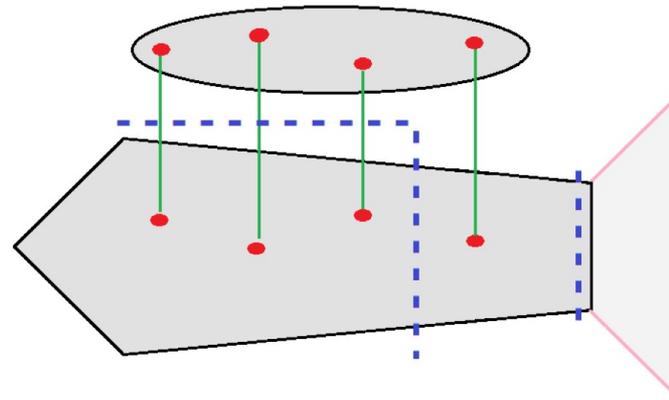
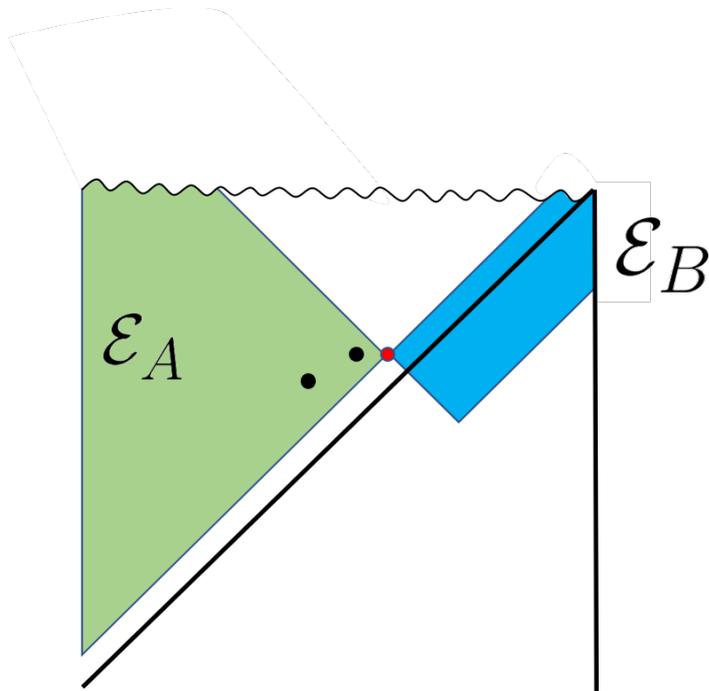


- RB radiation state is **pseudorandom**
- Computationally bounded observer O cannot destruct encoding of A in RB system.
- Implies that decoding Hawking radiation required exponential complexity

$$\text{complexity} \sim e^S$$

[Kim, Tang, Preskill et. al 2020]

Covariant python's lunch maximal cut (reverse sweep)



Maximinimax prescription:

$$S_{max}^{(gen)} = S_{rad} + 2 \frac{A_{hor}}{4\hbar G}$$

Post-selecting a qutrit

$$U|s\rangle|1\rangle = \frac{|\alpha\rangle|1\rangle + |\beta\rangle|2\rangle + |\gamma\rangle|3\rangle + |\delta\rangle|4\rangle}{2}$$

$$V \equiv \mathbf{1} - 2|\tilde{\alpha}\rangle\langle\tilde{\alpha}| = U^\dagger \left(\mathbf{1} \otimes (\mathbf{1} - 2|1\rangle\langle 1|) \right) U. \quad \left| \begin{array}{l} |\tilde{\alpha}\rangle \equiv U^\dagger|\alpha\rangle|1\rangle \quad ; \quad |\tilde{\beta}\rangle \equiv U^\dagger|\beta\rangle|2\rangle \\ |\tilde{\gamma}\rangle \equiv U^\dagger|\gamma\rangle|3\rangle \quad ; \quad |\tilde{\delta}\rangle \equiv U^\dagger|\delta\rangle|4\rangle \end{array} \right.$$

$$U_s \equiv 2|s\rangle\langle s| \otimes |1\rangle\langle 1| - \mathbf{1}$$

$$V|s\rangle|1\rangle = \frac{1}{2}(-|\tilde{\alpha}\rangle + |\tilde{\beta}\rangle + |\tilde{\gamma}\rangle + |\tilde{\delta}\rangle)$$

$$U_s V|s\rangle|1\rangle = |\tilde{\alpha}\rangle$$

$$U \left[U_s V \right] |s\rangle|1\rangle = |\alpha\rangle|1\rangle$$

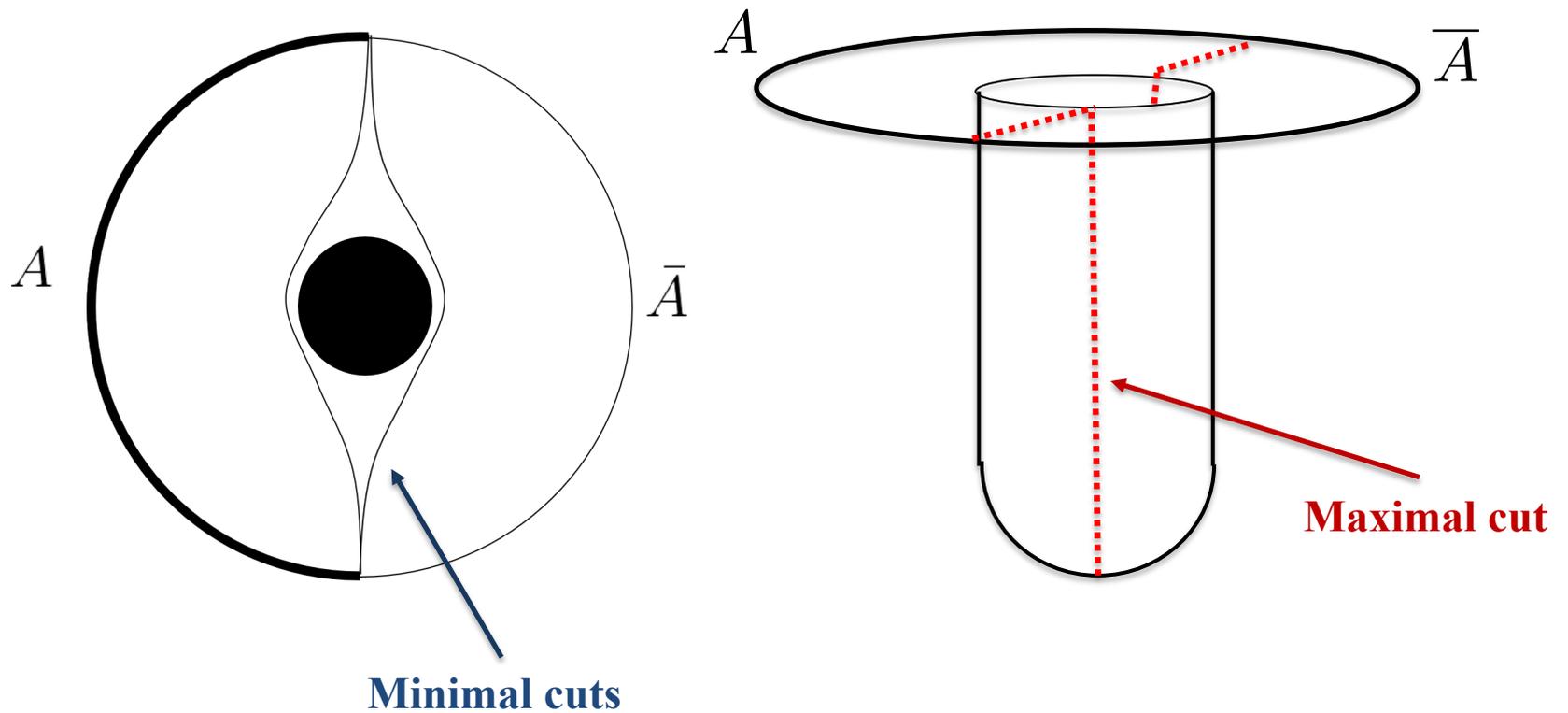
Post-selecting multiple qubits

$$U|s\rangle|0\rangle^{\otimes n} = \sin\theta|\alpha\rangle|0\rangle^{\otimes m} + \cos\theta|\beta\rangle$$

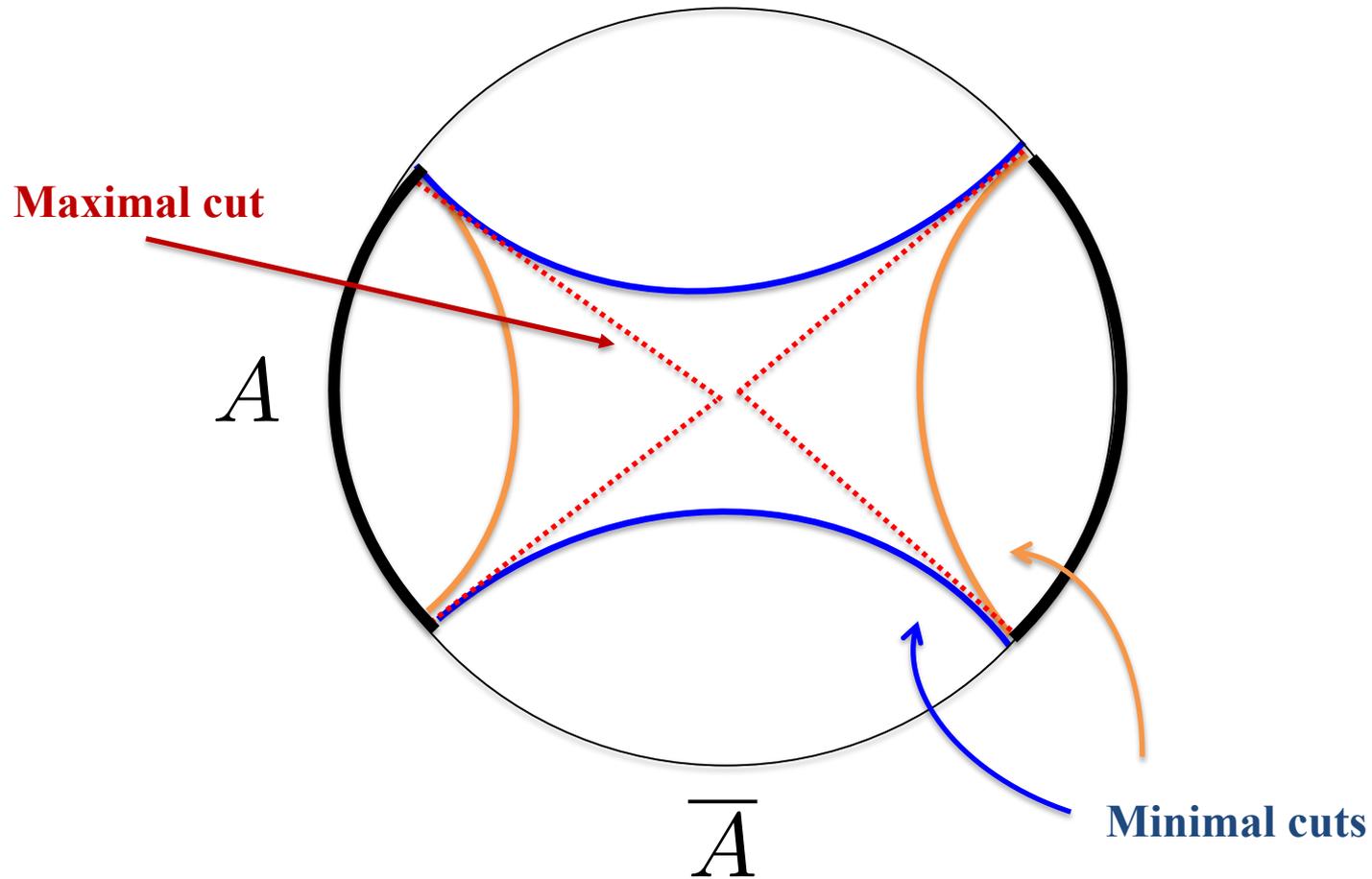
$$\begin{aligned} U_s &= 2|s\rangle\langle s| \otimes |0\rangle\langle 0|^{\otimes n} - \mathbf{1} \\ V &= U^\dagger \left(\mathbf{1} \otimes (\mathbf{1} - 2|0\rangle\langle 0|^{\otimes m}) \right) U \end{aligned}$$

$$U \left[U_s V \right]^l |s\rangle|0\rangle^{\otimes n} = \sin[(2l+1)\theta]|\alpha\rangle|0\rangle^{\otimes m} + \cos[(2l+1)\theta]|\beta\rangle$$

Python's lunch in pure state black hole

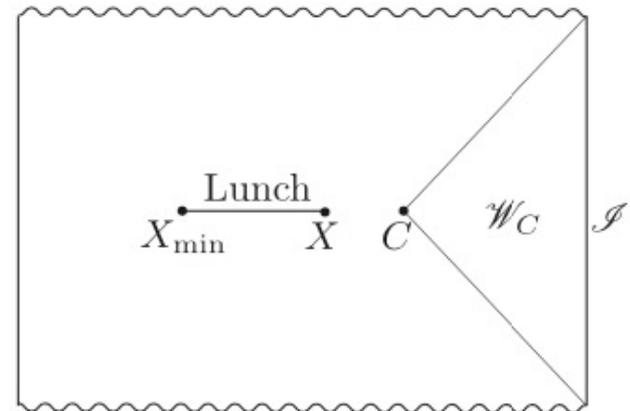
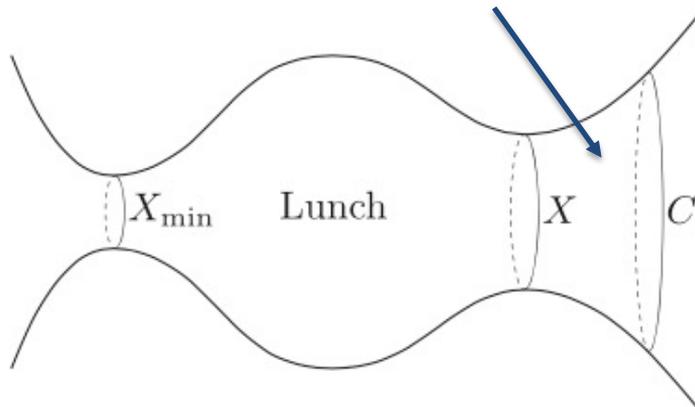


Python's lunch without black holes – empty AdS



Converse of Python's lunch conjecture

Easy to reconstruct



- HKLL construction
- Causal bulk propagation of perturbed boundary condition on Lorenzian timefolds

[Engelhardt, Penington, Shahbazi-Moghaddam 2021]

What's the complexity of removing post-selection / final state projection?

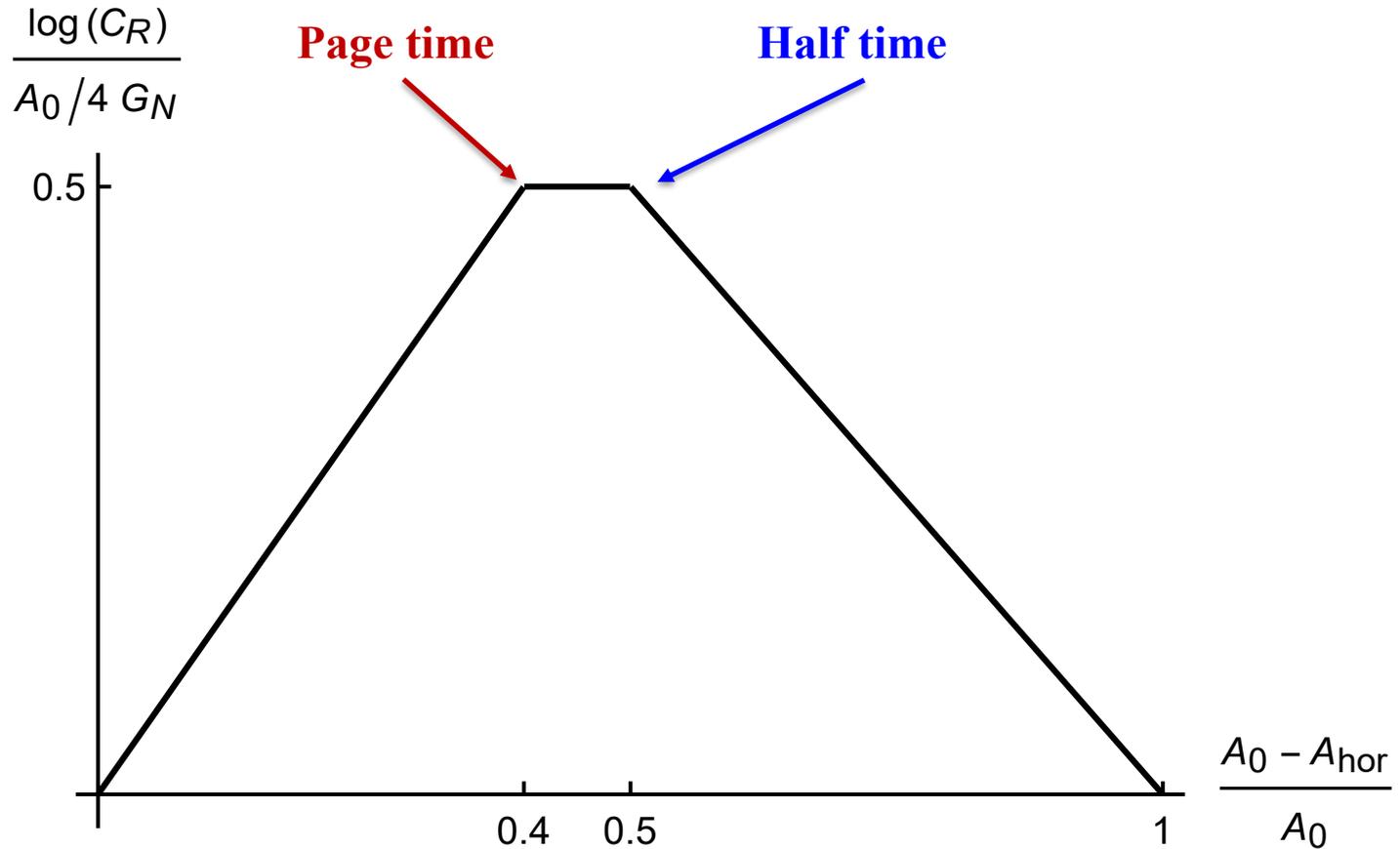
We can measure the qubits and hope to get the right answers. Probability of the right answer is $1/2^m$ and expected complexity

$$C \propto 2^{m_R} \cdot C_{TN}$$

It is not a fixed unitary and complexity can be improved.

We can do a generalized **Grover-like search**.

Restricted complexity for evaporating black hole



When are restricted and unrestricted complexities very different?

python's lunch



Final remarks

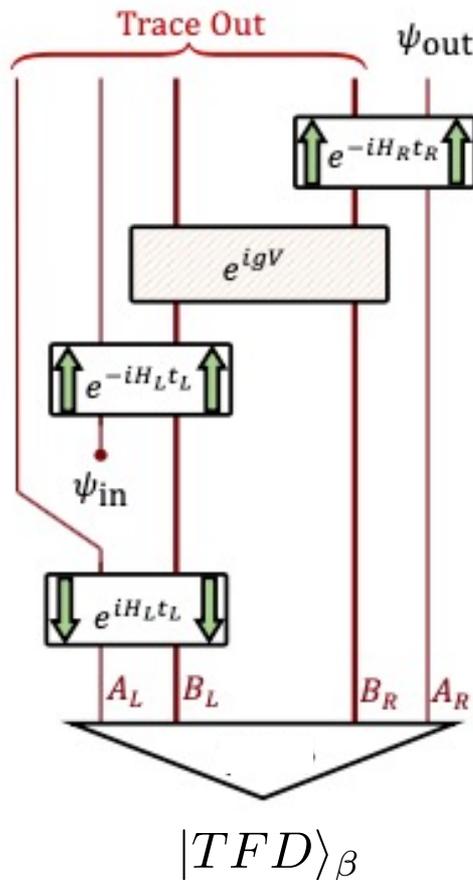
Summary: The restricted complexity of python's lunch is

$$C_R \propto C_{TN} \cdot \exp \left[\frac{1}{2} \left(S_{max}^{(gen)} - S_{min}^{(gen)} \right) \right]$$

Future directions

- Prove this formula beyond tensor network formalism, perhaps new definition of complexity is needed.
- Study other examples when python's lunches and non-minimal quantum extremal surfaces appear; pure state black holes and empty AdS.
 - What does this imply about reconstruction of operators in the python's lunch?
 - Relationship of python's lunch to quantum extended Church-Turing thesis

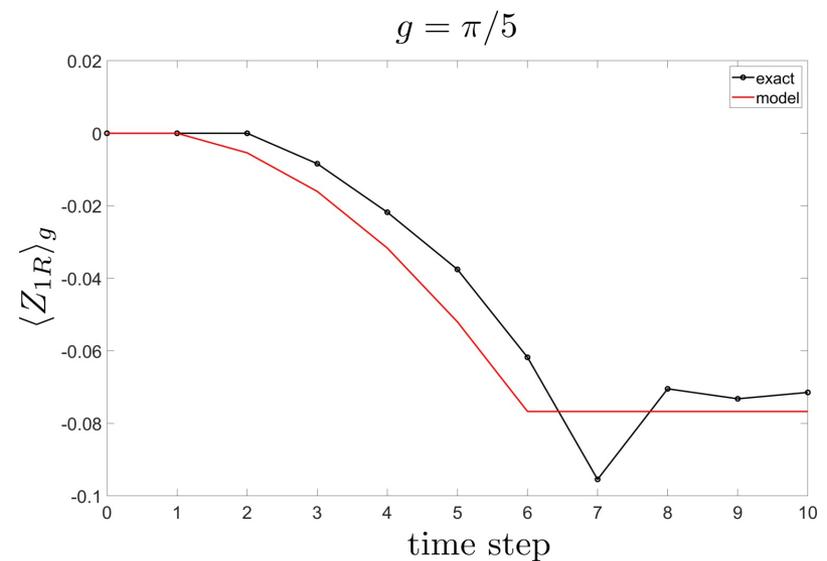
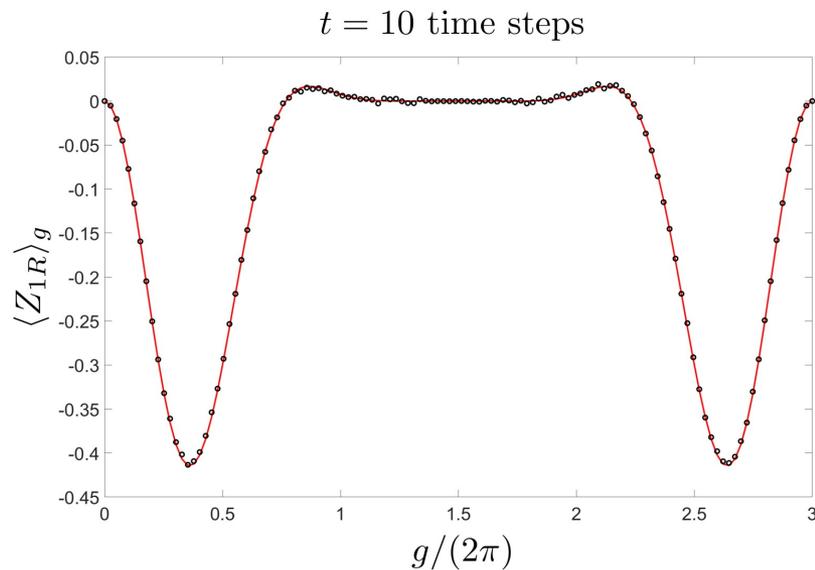
Some Interesting Questions



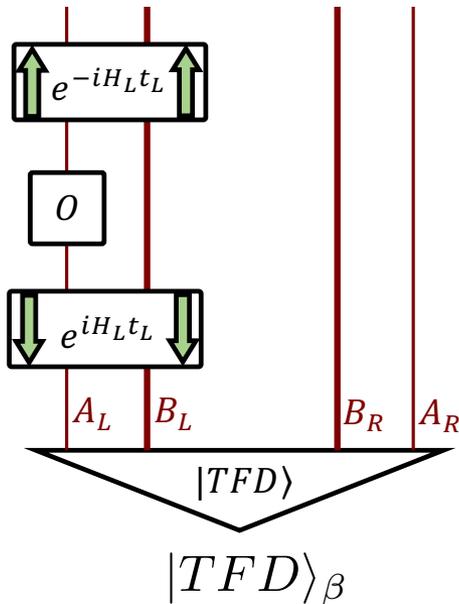
1. What are the Hamiltonians that can teleport quantum information?
2. Does the teleportation work at different temperatures?
3. How does the teleportation depend on the details of the coupling?

Regime #1: numerical study of the experimental proposal

Numerical evidence for Floquet version of Rydberg dynamics for N=7 qubits.



Regime #2: size winding explained



Operator growth

$$O(t) = e^{-iHt} O e^{+iHt} = \sum_P c_P(t) P$$

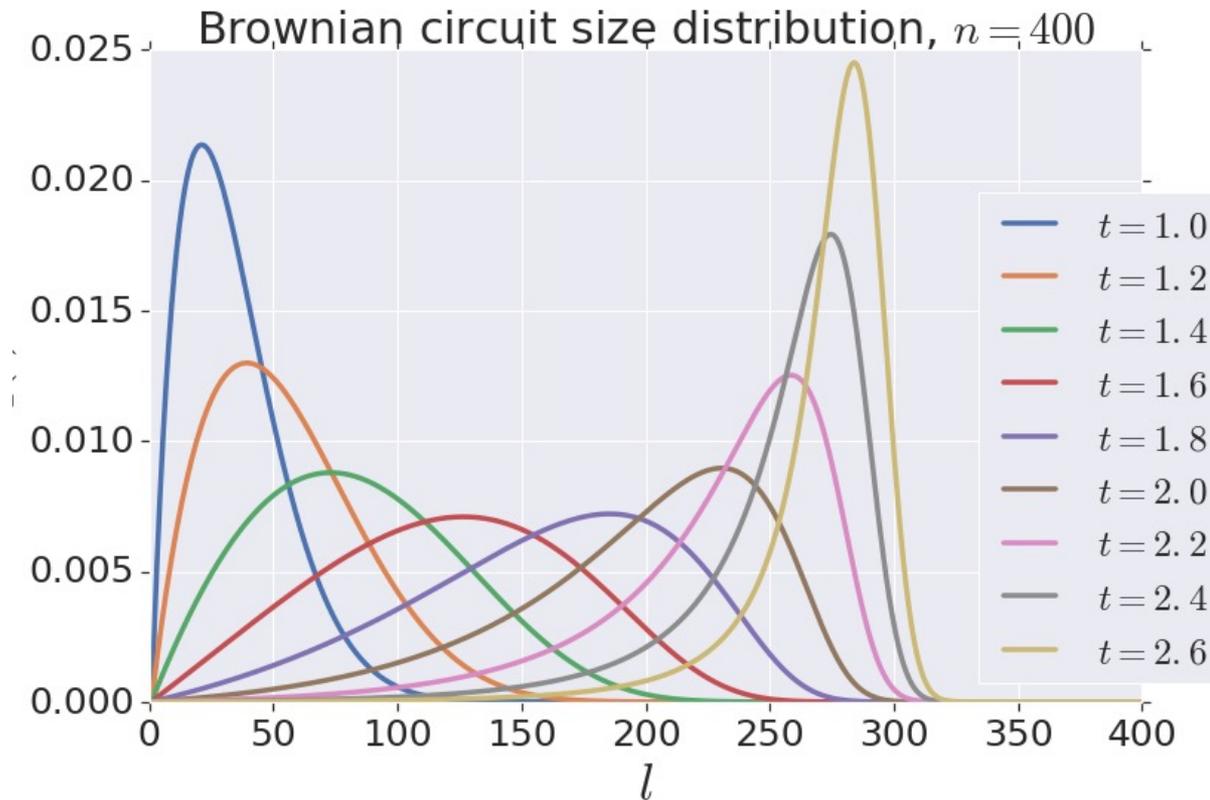
real

Operator size-distribution

$$p(l) = \sum_{|P|=l} |c_P|^2$$

For Hermitian operator and finite temperature coefficients are real.

Regime #2: Operator Size Distribution



Regime #2: size winding in SYK model

Size-winding is hard to illustrate from the boundary **analytically**.

Sachdev-Ye-Kitaev: $H_{SYK} = \sum_{i,j,k,l} J_{ijkl} \chi_i \chi_j \chi_k \chi_l$

Majorana fermions



We verify the results noticed earlier in JT gravity

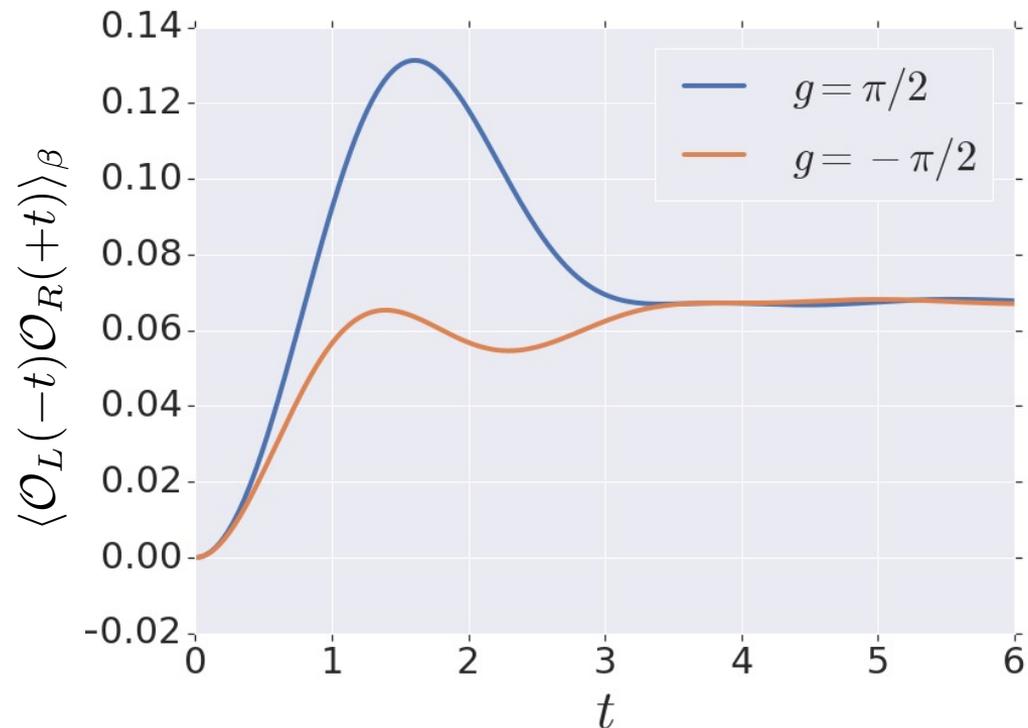
Moment = Size

Bulk AdS

Boundary CFT

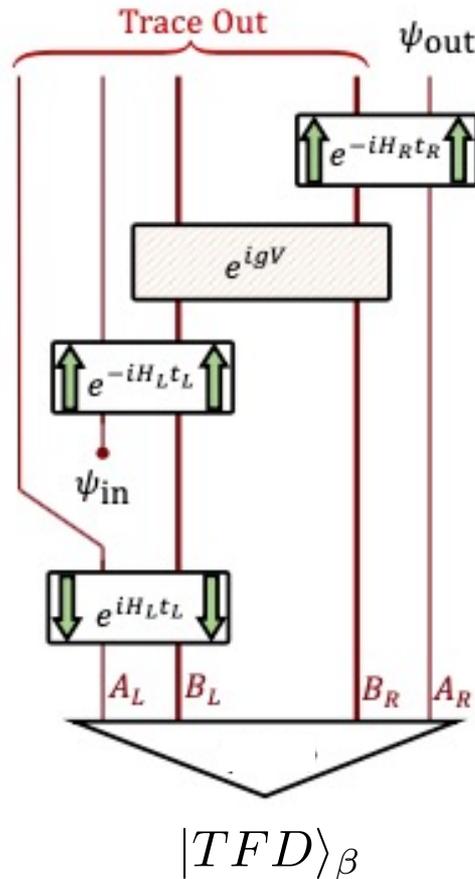
Regime #2: size winding beyond SYK model

Random Matrix Theory (GOE or GUE ensemble)



A weaker form of size-winding is responsible for capacity enhancement at low temperatures.

Summary



Phase One:

1. Infinite temperature (Bell pairs)
2. Any chaotic Hamiltonian
3. Single qubit teleportation

Phase Two:

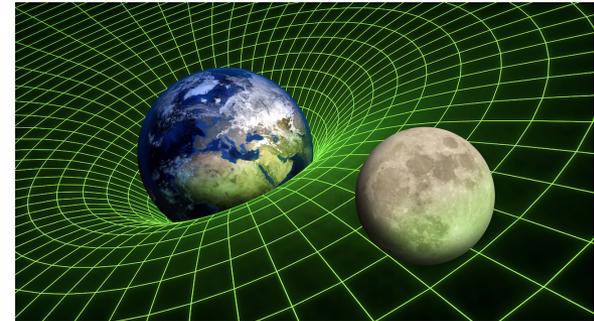
1. Finite temperature (TFD)
2. Hamiltonians that potentially have holographic duals
3. Many qubit teleportation

(Information travels through the wormhole)

Final remarks: quantum gravity in the lab



Holography



Quantum Devices:

- Rydberg atoms,
- Trapped ions,
- Superconducting qubits
- And others.....

Quantum Gravity:

- New holographic models
- Go beyond analytic large-N methods with exp.
- Stringy effects

Contributions

What can we learn about quantum gravity (holography) using the tools of quantum information and computation?

- ❖ Extend complexity = volume to evaporating black holes
- ❖ The role of python's lunch feature in the geometry to complexity
- ❖ Investigate other instances of python's lunch geometries

Can we study deep questions about quantum gravity in the near term quantum devices? (quantum gravity in the lab)

- ❖ Traversable wormholes, teleportation by size, and size winding
- ❖ Simulating quantum matrix models on a quantum computer
- ❖ Discovering new holographic models, protocols and observables